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Ann. Probab., **28 (2000), 1–35.**

Thick Points for Spatial Brownian Motion: Multifractal Analysis of Occupation Measure

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Let $\mathcal{T}(x,r)$ denote the total occupation measure of the ball of radius r centered at x for Brownian motion in $\mathbb{R}^3$. We prove that $\sup_{|x|\leq 1} \mathcal{T}(x,r)/(r^2|\log r|) \to 16/\pi^2$ a.s. as $r \to 0$, thus solving a problem posed by Taylor in 1974. Furthermore, for any $a \in (0,16/\pi^2)$, the Hausdorff dimension of the set of "thick points" x for which $\limsup_{r\to 0} \mathcal{T}(x,r)/(r^2|\log r|) = a$, is almost surely $2 - a\pi^2/8$; this is the correct scaling to obtain a nondegenerate "multifractal spectrum" for Brownian occupation measure. Analogous results hold for Brownian motion in any dimension $d > 3$. These results are related to the LIL of Ciesielski and Taylor (1962) for the Brownian occupation measure of small balls, in the same way that Lévy's uniform modulus of continuity,

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[2]Research partially supported by NSF grant #DMS-9704552 and by a US-Israel BSF grant.

[3]Research partially supported by NSF grants #DMS-9404391 and #DMS-9803597

[4]Research supported, in part, by grants from the NSF and from PSC-CUNY.

This research was supported, in part, by MSRI and by a US-Israel BSF grant.

AMS 1991 *subject classifications.* Primary 60J65, 28A80 ; secondary 60J55, 60F99, 28A78

Key words and phrases. thick points, multifractal analysis, occupation measure

and the formula of Orey and Taylor (1974) for the dimension of "fast points", are related to the usual LIL. We also show that the lim inf scaling of $\mathcal{T}(x,r)$ is quite different: we exhibit non-random $c_1, c_2 > 0$, such that $c_1 < \sup_x \liminf_{r\to 0} \mathcal{T}(x,r)/r^2 < c_2$ a.s. In the course of our work we provide a general framework for obtaining lower bounds on the Hausdorff dimension of random fractals of 'limsup type'.

1. Introduction. For any Borel measurable function f from $0 \le t \le T$ to $\mathbb{R}^d$ we denote by μ_T^f its *occupation measure*:

$$\mu_T^f(A) = \int_0^T \mathbf{1}_A(f_t)\,dt$$

for all Borel sets $A \subseteq \mathbb{R}^d$. Throughout, $B(x,r)$ denotes the ball in $\mathbb{R}^d$ of radius r centered at x, and $\{W_t\}_{t\ge 0}$ denotes Brownian motion in $\mathbb{R}^d$, $d \ge 3$.

In the last decade, much insight into the structure of various measures has been gained from their *multifractal analysis.* A general introduction to this analysis can be found in Olsen [12], Riedi [21] and Falconer [8]; certain important random measures were analyzed by Hu-Taylor [9], Taylor [28], Perkins-Taylor [20], Lawler [11] and Shieh-Taylor [24].

Consider Brownian occupation measure μ_T^W in $\mathbb{R}^d$, $d \ge 3$. It is well known that for almost all Brownian paths W, the pointwise Hölder exponent

$$\mathsf{H\ddot{o}lder}(\mu_T^W, x) := \lim_{\epsilon\to 0} \frac{\log \mu_T^W(B(x,\epsilon))}{\log \epsilon} \tag{1.1}$$

takes the value 2 for *all* points x in the range $\{W_t \,\big|\, 0 \le t \le T\}$. In particular, the usual multifractal spectrum $a \mapsto \dim\{x \in \mathbb{R}^d : \mathsf{H\ddot{o}lder}(\mu_T^W, x) = a\}$ vanishes for all $a \ne 2$, $a > 0$. Indeed, this fact was crucial in Kaufman's work [10], written long before the term "multifractal" was invented.

Rather than being the end of the story, this means that standard multifractal analysis must be refined to capture the delicate fluctuations of occupation measure under scaling; the problem of obtaining such a refined analysis was posed by Hu and Taylor [9, Pg. 287] in 1997, but it is closely linked to problems posed by Taylor [26] in 1974. Our main results, Theorems 1.1 and 1.3 below, resolve these problems.

The correct scaling for studying the fluctuations of occupation measure was already indicated by Taylor [26]; more details were given by Perkins-Taylor [18, Lemmas 2.3 and 2.5], who showed that there exist absolute constants $0 < c_1 < c_2 < \infty$, such that almost surely for all points $x \in \{W_t | \, 0 \leq t \leq T\}$ and all positive $\epsilon \leq \epsilon_0(\omega)$,

$$c_1 \epsilon^2 / |\log \epsilon| \leq \mu_T^W(B(x,\epsilon)) \leq c_2 \epsilon^2 |\log \epsilon| \,. \tag{1.2}$$

(As they point out, the lower bound is immediate from Lévy's uniform modulus of continuity.)

Our main result describes the multifractal nature, in a fine scale, of "thick points" for the occupation measure of Brownian motion in $I\!R^d$, $d \geq 3$. (We call a point $x \in I\!R^d$ on the Brownian path a *thick point* if x is in the set considered in (1.3) for some $a > 0$; similiarly, $t > 0$ is called a $\mathsf{thick\ time}$ if it is in the set Thick_a considered in (1.4) for some $a > 0$ and $T > 0$.)

THEOREM 1.1. *With $d \geq 3$, let q_d denote the first positive zero of the Bessel function $J_{d/2-2}(x)$. (See [31] for information on q_d; in particular, $q_3 = \pi/2$.) Then, for any $T \in (0, \infty]$ and all $0 < a \leq 4/q_d^2$,*

$$\dim\{x \in I\!R^d \,\Big|\, \limsup_{\epsilon \to 0} \frac{\mu_T^W(B(x,\epsilon))}{\epsilon^2 |\log \epsilon|} = a\} = 2 - a q_d^2/2 \quad \text{a.s.} \tag{1.3}$$

Equivalently, for any $T \in (0,\infty]$ and all $0 < a \le 4/q_d^2$,

$$\dim\{0 \le t < T \,\Big|\, \limsup_{\epsilon\to 0} \frac{\mu_T^W(B(W_t,\epsilon))}{\epsilon^2|\log\epsilon|} = a\} = 1 - aq_d^2/4 \quad \text{a.s.} \tag{1.4}$$

Denote the set in (1.4) by Thick_a*. Then* $\mathsf{Thick}_a \neq ?$ *at the critical value* $a = 4/q_d^2$*.*

For comparison purposes, recall three fundamental results on Brownian increments:

(i) The large increments at a fixed time t, are governed by Khinchin's classical LIL:

$$\limsup_{\epsilon\to 0} \frac{W_{t+\epsilon} - W_t}{(2\epsilon \log|\log\epsilon|)^{1/2}} = 1 \quad \text{a.s.}$$

(ii) The dimension of certain exceptional *fast points* was determined by Orey-Taylor [13]:

$$\forall a \in [0,1], \quad \dim\left\{0 \le t < T \,\Big|\, \limsup_{\epsilon\to 0} \frac{W_{t+\epsilon} - W_t}{(2\epsilon|\log\epsilon|)^{1/2}} = a\right\} = 1 - a^2 \quad \text{a.s.}$$

(This can be viewed as a multifractal decomposition of white noise.)

(iii) Lévy's uniform modulus of continuity governs the largest increments overall:

$$\lim_{\epsilon\to 0} \sup_{0\le t\le T} \frac{W_{t+\epsilon} - W_t}{(2\epsilon|\log\epsilon|)^{1/2}} = 1 \quad \text{a.s.}$$

The three statements above hold in any dimension $d \ge 1$. Next, we indicate their analogues for Brownian occupation measure in dimension $d \ge 3$; only the first of these was previously known.

(i') The limsup asymptotic behavior of Brownian occupation measure around a fixed time t, is governed by the LIL of Ciesielski-Taylor [3, Theorem 3]:

for any $T \in (0, \infty]$ and $t \leq T$,

$$\limsup_{\epsilon \to 0} \frac{\mu_T^W(B(W_t, \epsilon))}{\epsilon^2 \log|\log \epsilon|} = \frac{2}{q_d^2} \quad \text{a.s.} \tag{1.5}$$

(ii') The dimension of exceptional thick times is given by (1.4) above.

(iii') Our results (1.7) and (1.9) give the largest occupation measure possible for a small ball.

Further remarks on Theorem 1.1.

- Perhaps more significant than the numerical values obtained in (1.3) and (1.4) is the insight gained, while proving these results, about the manner by which the "thick points" on the Brownian path arise. The key to our proof of Theorem 1.1 is a **localization phenomenon** for transient Brownian motion: the balls of radius ϵ that have the largest occupation measure (of order $\epsilon^2|\log \epsilon|$), accumulate most of this measure in a surprisingly short time interval (of length at most $\epsilon^2|\log \epsilon|^b$ for some b, e.g. $b = 6$ works); see Section 3 where this localization is established. The localization phenomenon breaks down in dimension $d = 2$, where the correct scaling of occupation measure, and the techniques needed to establish it, are quite different. In [6] we have obtained the corresponding results for the planar case; we emphasize that the current paper concerns only $d \geq 3$.

- Given the localization phenomenon, there are several possible approaches to the proof of the lower bound in (1.4). Our proof relies on a general lower bound on Hausdorff measure of random fractals "of limsup type", Theorem 2.1. This general bound sharpens similar estimates obtained by

Orey-Taylor [13], Hu-Taylor [9], Deheuvels-Mason [4] and Shieh-Taylor [24]; of course, our work owes a substantial debt to these earlier papers.

- For any $x \notin \{W_t \,\big|\, 0 \le t \le T\}$ and ϵ small enough, $\mu_T^W(B(x,\epsilon)) = 0$. Hence, the equivalence of (1.3) and (1.4) is a direct consequence of the *uniform dimension doubling* property of Brownian motion, due to Kaufman [10] (see also, [18, Eqn. (0.1)]).

- Let v_{d-2} denote the first eigenvalue of $(1/2)\Delta$ in the unit ball of $I\!R^{d-2}$ with Dirichlet boundary conditions. As the spherically symmetric fundamental solution for the Laplacian eigenvalue problem in $B(0,1)$ is $J_{d/2-2}(\sqrt{\lambda}|x|)$, the required Dirichlet boundary conditions imply that $v_{d-2} = q_d^2/2$ (see for example [3, (2.15)]). The appearance of $(d-2)$ in our result for d dimensions is due to the celebrated identity of Ciesielski-Taylor [3, Theorem 2].

To indicate the qualitative difference between the sets of thick points and the most familiar random fractals associated with Brownian motion (the range, the graph, and the level sets) we present the following proposition; for the definition and properties of packing dimension $\dim_{\mathrm{P}}$, see [29] or [8].

Proposition 1.2. *Let the notation of Theorem 1.1 be in force. For all $0 < a \le 4/q_d^2$, the union $\mathsf{Thick}_{\ge a} := \cup_{b \ge a} \mathsf{Thick}_b$ has the same Hausdorff dimension as Thick_a a.s., but its packing dimension a.s. satisfies $\dim_{\mathrm{P}}(\mathsf{Thick}_{\ge a}) = 1$. Equivalently,*

$$\dim_{\mathrm{P}}\{x \in I\!R^d \,\Big|\, \limsup_{\epsilon \to 0} \frac{\mu_T^W(B(x,\epsilon))}{\epsilon^2 |\log \epsilon|} \ge a\} = 2 \quad \text{a.s.} \tag{1.6}$$

Remark. The importance of comparing the Hausdorff and packing dimensions of a set was stressed in the survey Taylor [27]. By a more involved argument, it can be shown that Thick_a itself also has packing dimension 1 for $0 < a \leq 4/q_d^2$. (For $a = 4/q_d^2$, this statement follows from Proposition 1.2.)

The next theorem solves two problems posed by Taylor in 1974 (see [26, Pg. 201]).

THEOREM 1.3. *Let $\{W_t\}$ be a Brownian motion in $\mathbb{R}^d$, $d \geq 3$. Then, for any $R \in (0, \infty)$ and any $T \in (0, \infty]$,*

$$\lim_{\epsilon \to 0} \sup_{|x| \leq R} \frac{\mu_T^W(B(x,\epsilon))}{\epsilon^2 |\log \epsilon|} = 4q_d^{-2} \quad \text{a.s.} \tag{1.7}$$

Furthermore, for any $k \in (0, \infty)$ and any $T \in [k, \infty]$,

$$\lim_{\epsilon \to 0} \inf_{t \in [0,k]} \frac{\mu_T^W(B(W_t,\epsilon))}{\epsilon^2 / |\log \epsilon|} = 1 \quad \text{a.s.} \tag{1.8}$$

Remarks:

- Our proof shows that for any $T \in (0, \infty]$,

$$\lim_{\epsilon \to 0} \sup_{0 \leq t \leq T} \frac{\mu_T^W(B(W_t,\epsilon))}{\epsilon^2 |\log \epsilon|} = 4q_d^{-2} \quad \text{a.s.} \tag{1.9}$$

- Combining (1.3) and (1.7) we see that

$$\sup_{x \in \mathbb{R}^d} \limsup_{\epsilon \to 0} \frac{\mu_\infty^W(B(x,\epsilon))}{\epsilon^2 |\log \epsilon|} = 4q_d^{-2} \quad \text{a.s.}$$

  In particular, the sets in (1.3) and (1.4) are a.s. empty for any $a > 4q_d^{-2}$, $T \in (0, \infty]$.

- A detailed multifractal analysis of thin times for Brownian motion, that is $t \in [0, T]$ satisfying

$$\liminf_{\epsilon \to 0} \frac{\mu_T^W(B(W_t,\epsilon))}{\epsilon^2 / |\log \epsilon|} = a$$

  for some $a \geq 1$, can be found in [7]. The relation between thin times and (1.8) is the same as the relation between thick times and (1.7).

Computation of Laplace transforms is an important component of a complete multifractal analysis, and it was also the starting point of our investigation. Pemantle, Peres and Shapiro [15] showed that $\int_0^1 \mu_1^W(B(W_t,\epsilon))/\epsilon^2\,dt$, the pathwise first moment of the ratio $\mu_1^W(B(W_t,\epsilon))/\epsilon^2$, remains bounded almost surely as $\epsilon\to 0$. The following theorem provides a pathwise asymptotic formula for the moment generating function of that ratio. In one sense, it is finer than Theorem 1.1, since it yields a precise estimate of the total duration in $[0,1]$ that the Brownian particle spends in balls of radius ϵ that have unusually high occupation measure (see Corollary 1.5 below). Such an estimate (which is an analogue in our setting of the "coarse multifractal spectrum", cf. Riedi [21]), cannot be inferred from Theorem 1.1.

THEOREM 1.4. *Denote by $\bar\mu_\infty^{\bar W}$ the total occupation measure for a two-sided Brownian motion $\{\bar W_t\}_{-\infty}^{\infty}$ in $\mathbb{R}^d$, $d\ge 3$. Then for each $\theta<q_d^2/2$,*

$$\lim_{\epsilon\to 0}\int_0^1 e^{\theta\mu_1^W(B(W_t,\epsilon))/\epsilon^2}\,dt=\mathbb{E}\left(e^{\theta\bar\mu_\infty^{\bar W}(B(0,1))}\right)\qquad a.s. \tag{1.10}$$

Remarks:

- We note by [3] that

$$\mathbb{E}\left(e^{\theta\bar\mu_\infty^{\bar W}(B(0,1))}\right)=\left(\mathbb{E}\left(e^{\theta\mu_\infty^{W}(B(0,1))}\right)\right)^2=\frac{1}{\prod_{j=1}^\infty\left(1-\frac{2\theta}{q_{d,j}^2}\right)^2} \tag{1.11}$$

for each $\theta<q_d^2/2$, where $\{q_{d,j}\}_{j\ge 1}$ are the positive zeros of the Bessel function $J_{d/2-2}(x)$, enumerated in increasing order. It is clear that the right hand side diverges as $\theta\uparrow q_d^2/2=q_{d,1}^2/2$. The case $d=3$ is particularly explicit because then $q_3=\pi/2$ and the right hand side of (1.11) simplifies to $\cos^{-2}(\sqrt{2\theta})$ (c.f. [3]).

- Let τ denote a random variable uniform on $[0,1]$, which is independent of the Brownian path W. Then, (1.10) implies in particular that for almost every Brownian path W, the ratio $\mu_1^W(B(W_\tau,\epsilon))/\epsilon^2$, a random variable in τ, converges in law as $\epsilon \to 0$ to the total occupation time $\bar{\mu}_\infty^{\bar{W}}(B(0,1))$ of the unit ball by a two-sided Brownian motion $\bar{W}$.

Next, we state the promised corollary of Theorem 1.4, which is analogous to the coarse multifractal spectrum.

Corollary 1.5. *Let $\{W_t\}$ be a Brownian motion in $\mathbb{R}^d$, $d \geq 3$, and denote Lebesgue measure on $\mathbb{R}^1$ by $\mathcal{L}eb$. Then, for any $a \in (0, 4/q_d^2)$,*

$$\lim_{\epsilon\to 0} \frac{\log \mathcal{L}eb\Big\{0 \leq t \leq 1 \,\Big|\, \mu_1^W(B(W_t,\epsilon)) \geq a\epsilon^2 |\log \epsilon|\Big\}}{\log \epsilon} = aq_d^2/2 \quad \text{a.s.}$$

The thick points considered in Theorem 1.1 are centers of balls $B(x,\epsilon)$ with unusually large occupation measure for infinitely many radii, but these radii might be quite rare. The next theorem shows that for the balls $B(x,\epsilon)$ to have unusually large occupation measure for *all* small radii ϵ and the same center x, what constitutes "unusually large" must be interpreted more modestly. Define

$$I_d(a) := \frac{a}{4}(\max\{0, d-2-\frac{2}{a}\})^2 \,, \tag{1.12}$$

and let

$$C_d := \inf\{a : I_d(a) = 2\} = \frac{2}{d - 2\sqrt{d-1}} \,. \tag{1.13}$$

(The equality on the right is easily verified.)

Then

Theorem 1.6. *For $\{W_t\}$ a Brownian motion in $\mathbb{R}^d$, $d \geq 3$, and $a \in (0, C_d]$,*

$$\dim\{x \in \mathbb{R}^d \,\Big|\, \liminf_{\epsilon\to 0} \frac{\mu_\infty^W(B(x,\epsilon))}{\epsilon^2} \geq a\} \leq 2 - I_d(a) \quad \text{a.s.} \tag{1.14}$$

and this can be strengthened to

$$\dim_{\mathrm{P}}\{x \in \mathbb{R}^d \;\Big|\; \liminf_{\epsilon\to 0} \frac{\mu_\infty^W(B(x,\epsilon))}{\epsilon^2} \geq a\} \leq 2 - I_d(a) \quad \text{a.s.} \tag{1.15}$$

where $\dim_{\mathrm{P}}$ *denotes packing dimension. Moreover,*

$$\frac{1}{d} \leq \sup_{x\in\mathbb{R}^d} \liminf_{\epsilon\to 0} \frac{\mu_\infty^W(B(x,\epsilon))}{\epsilon^2} \leq C_d \quad \text{a.s.} \tag{1.16}$$

Remarks:

- In particular, replacing the lim sup by lim inf in (1.3) and (1.4) yields an a.s. empty set for any $a > 0$.
- The new assertion in (1.16) is the right hand inequality; the inequality on the left is an immediate consequence of Theorem 9 of Perkins [17] concerning "Brownian slow points".
- It is an open problem to determine exactly the dimension appearing in (1.14) and the precise asymptotics in (1.16).
- That the upper bound (1.14) on Hausdorff dimension applies to packing dimension as well is in sharp contrast with Theorem 1.1 and Proposition 1.2. Intuitively, the reason for this contrast is that for a point to be in the set considered in (1.3), it only needs to satisfy a certain condition at infinitely many scales, so that set can appear large at other scales; these scales can be used to pack many disjoint balls with centers in the set. Points considered in (1.14), however, must satisfy a (less stringent) condition at *all* scales.

The next section contains a discussion of fractals "of limsup type" and a general lower bound (Theorem 2.1) for their Hausdorff measure. In Section 3 we

prove the crucial Localization Lemma 3.1. The results of those two sections are applied in Section 4 to establish the lower bounds on Hausdorff dimension in Theorem 1.1 and Proposition 1.2. The complementary upper bounds in Theorem 1.1 are proved in Section 5. Combining these bounds with the Localization Lemma 3.1, we prove Theorem 1.3 in Section 6. Section 7 is devoted to the proof of Theorem 1.4, with Corollary 1.5 proved in Section 8. Theorem 1.6 is proved in Section 9. At the end of the paper we present some open problems.
Analogous results for transient symmetric stable processes are contained in [5].

2. Random fractals of limsup type. Suppose that for each $n \geq 1$, a finite union $A(n)$ of intervals of length λ_n is given. Assume that $\lambda_n \to 0$ as $n \to \infty$, and that the number of intervals comprising $A(n)$ grows like a negative power of λ_n. We call $A := \limsup A(n) = \bigcap_{n=1}^{\infty} \bigcup_{k=n}^{\infty} A(k)$ a *fractal of limsup type*. We will be interested in situations in which the $A(n)$ are random, and in hypotheses on their distribution which will allow us to obtain dimension bounds on A. The main result of this section, Theorem 2.1, provides a general framework for obtaining lower bounds on the Hausdorff measure of random fractals of limsup type.

Random sets that are (well approximated by) random fractals of limsup type include:

- The *fast points* of Orey-Taylor [13];
- The initial points of exceptional Brownian excursions considered by Barlow-Perkins [1];
- The *close approaches* on the Brownian path measured by Perkins-Taylor [19];

- The paths in a family tree where a tree-indexed random walk has positive *burst speed*, see Benjamini-Peres [2];

- Times where the Strassen functional LIL fails, see Deheuvels-Mason [4];

- Sets arising in multifractal analysis of stable subordinators (studied by Hu-Taylor [9] and by Shieh-Taylor [24]).

- The sets Thick_a in Theorem 1.1 .

Such random sets differ qualitatively from the random fractals most frequently encountered (e.g. ranges, graphs, levels sets and *slow points* of Brownian motion). For instance, the packing dimension of sets of limsup type is typically full, hence larger than their Hausdorff dimension; see Corollary 2.4. In particular, that corollary implies that the sets of fast points of [13] have packing dimension 1 (The assertion to the contrary in [27, Pg. 401] is wrong).

Three general methods have been employed to establish lower bounds for Hausdorff dimension of random fractals of limsup type. (These methods were used earlier for other sets).

- Orey-Taylor [13] constructed a Frostman measure directly, using estimates on binomial probabilities. Their method is expounded by Deheuvels-Mason [4]. This elegant method requires strong independence assumptions "within levels", and it is difficult to refine it to handle sets defined by an equality, like Thick_a, rather than an inequality. Orey-Taylor [13, Pg. 185] state that this can be done for the random fractals of limsup type which they consider, the Brownian fast points, by "tightening their argument", but extending this to more general situations seems quite hard.

- Intersection properties with an independent random set (the range of a stable subordinator) were used by Barlow-Perkins [1] and Perkins-Taylor [19]; random Cantor sets arising from fractal percolation as in [16] could also be used. Here independence assumptions can be replaced by correlation bounds, but, as above, handling sets like Thick_a is unwieldy.

- A powerful method based on estimation of energy integrals was used by Hu-Taylor [9] and Shieh-Taylor [24]. Below we sharpen and extend this method, and show that it yields good estimates of Hausdorff measure, while requiring only mild correlation hypothesis.

Let $\mathcal{D}_n$ denote the collection of dyadic intervals $\{[(i-1)2^{-n}, i2^{-n}]\}_{i=1}^{2^n}$. For any increasing function $\varphi : [0,1] \to [0,\infty)$ with $\varphi(0) = 0$, let $\mathcal{H}^\varphi(A)$ denote the Hausdorff measure of a set A in the gauge φ (see, e.g., [27] for the definition).

Theorem 2.1. *Suppose that for every $n \geq 1$, a collection of $\{0,1\}$ valued random variables $\{Z_I\}_{I\in\mathcal{D}_n}$ is given, so that $p_n := \mathbf{P}(Z_I = 1)$ is the same for all $I \in \mathcal{D}_n$. Let*

$$A(n) = \cup\{I \in \mathcal{D}_n \,\Big|\, Z_I = 1\} \text{ and } A := \limsup A(n) = \bigcap_{n=1}^{\infty}\bigcup_{k=n}^{\infty} A(k)\,.$$

For $I \in \mathcal{D}_m$, with $m < n$, define

$$M_n(I) = \sum_{J\in\mathcal{D}_n, J\subset I} Z_J\,.$$

Choose $\zeta(n) \geq 1$ such that

$$\mathrm{Var}(M_n(I)) \leq \zeta(n)\mathbb{E}(M_n(I)) = \zeta(n)p_n 2^{n-m}\,. \tag{2.1}$$

Let $\varphi(r)$ be a gauge function which is regularly varying of index $\alpha \in (0,1)$ as

$r \downarrow 0$. (I.e., $\varphi(r) = r^\alpha L(r)$ where $L(cr)/L(r) \to 1$ as $r \downarrow 0$ for any $c > 0$). If

$$\frac{2^{-n}\zeta(n)}{\varphi(2^{-n})p_n} \to 0 \tag{2.2}$$

then $\mathcal{H}^\varphi(A) > 0$ a.s.

Remarks:

- We emphasize that no independence or correlation assumptions are made relating Z_I and Z_J for I and J of different lengths.
- $\mathcal{H}^\varphi(A) > 0$ immediately implies that $\dim(A) \geq \alpha$.
- Theorem 2.1 can be applied to the "fast points" and "thick points" of a variety of processes; the only essential requirements are stationarity of increments, suitable decay of correlations and (for discontinuous processes) bounds on the jump probabilities. The non-vanishing of Hausdorff *measure* is proved in Theorem 2.1, rather than merely a bound on dimension, in order to handle the sets Thick_a, rather than just their unions $\mathsf{Thick}_{\geq a} = \cup_{b \geq a}\mathsf{Thick}_b$, in (1.4). (See the remark following the statement of Corollary 4.1).
- Let $\varphi_n := 1/\varphi(2^{-n})$. Beyond the obvious fact of some exponential growth of φ_n, our proof uses only the following simple consequence of the assumption that $\varphi(r)$ is regularly varying of index $\alpha \in (0,1)$ as $r \downarrow 0$: for some $C < \infty$ that does not depend on n,

$$\sum_{m=1}^{n} \varphi_m \leq C\varphi_n \qquad \text{and} \qquad \sum_{m=n}^{\infty} 2^{-m}\varphi_m \leq C2^{-n}\varphi_n. \tag{2.3}$$

- We will apply Theorem 2.1 below to prove Theorem 1.1. In that application, we will take $\varphi(r) = r^{1-\gamma}|\log_2(r)|^{13}$ with $p_n \geq 2^{-n\gamma}$ for some $0 < \gamma < 1$

and $\zeta(n) = n^{12}$, where throughout this paper, $\log_2$ stands for the logarithm to the base 2.

- Theorem 2.1, which is formulated for random fractals of limsup type in $[0,1]$, has an obvious generalization to random 'fractals of limsup type' in $[0,1]^d$. In this setup we can take $\varphi(r)$ to be any gauge function which is regularly varying of index $\alpha \in (0,d)$ as $r \downarrow 0$, and replace (2.1) and (2.2) by $\mathrm{Var}(M_n(I)) \le \zeta(n)\mathbb{E}(M_n(I)) = \zeta(n)p_n 2^{d(n-m)}$ and $2^{-dn}\zeta(n)/(\varphi(2^{-n})p_n) \to 0$ respectively. The proof of such a generalization is basically identical to the proof of Theorem 2.1.

To establish Theorem 2.1 we need two lemmas. The first one is a version of the well-known connection between energy and Hausdorff measure. For the reader's convenience, we include the brief proof.

LEMMA 2.2. *Fix an increasing gauge function φ such that $\varphi(0) = 0$. Suppose that B is a Borel set in $[0,1]$, and ν is a probability measure on B. If the dyadic energy*

$$\mathcal{E}_\varphi(\nu) := \sum_{m=1}^{\infty} \sum_{J \in \mathcal{D}_m} \frac{\nu(J)^2}{\varphi(2^{-m})}$$

of ν is finite, then $\mathcal{H}^\varphi(B) > 0$.

(In fact $\mathcal{H}^\varphi(B) = \infty$, but that is unimportant for our purpose). See [14] for the connection of $\mathcal{E}_\varphi(\nu)$ to more traditional expressions for energy.

Proof: Let

$$\Psi(x) := \sum_{m=1}^{\infty} \sum_{J \in \mathcal{D}_m} \frac{\nu(J)}{\varphi(2^{-m})} \mathbf{1}_J(x)\,.$$

Since $\int_B \Psi(x)\,d\nu = \mathcal{E}_\varphi(\nu)$, taking $C = 2\mathcal{E}_\varphi(\nu)$, the set $B_C := \{x \in B \,\big|\, \Psi(x) \le C\}$ has $\nu(B_C) \ge 1/2$. The restriction ν_C of ν to B_C satisfies $\nu_C(J) \le C\varphi(2^{-m})$ for every $J \in \mathcal{D}_m$ for all m. Since any interval $I \subset [0,1]$ can be covered by three shorter dyadic intervals, it follows that $\nu_C(I) \le 3C\varphi(|I|)$ for any interval I. Hence, if $\mathcal{A}$ is any countable collection of intervals with $B_C \subseteq \cup_{\mathcal{A}} I$, then

$$\frac{1}{2} \le \nu(B_C) \le \sum_{\mathcal{A}} \nu_C(I) \le 3C \sum_{\mathcal{A}} \varphi(|I|)$$

which implies that $1/(6C) \le \mathcal{H}^\varphi(B_C)$.

Alternatively, the a.s. finiteness of Ψ, in conjunction with [23], imply that $\mathcal{H}^\varphi(B) = \infty$. $\square$

The following lemma, which, roughly speaking, controls the "quadratic variation" of the random sets $A(n)$, is the key to the proof of Theorem 2.1. Recall that $\varphi_n = 1/\varphi(2^{-n})$, and note that by (2.2), for any ℓ we can choose an integer $n(\ell) > \ell$ such that

$$\frac{\varphi_{n(\ell)}\zeta(n(\ell))}{2^{n(\ell)} p_{n(\ell)}} \le 2^{-3\ell}. \tag{2.4}$$

Lemma 2.3. *Let the assumptions of Theorem 2.1 be in force. There exist an a.s. finite random variable $\ell_0(\omega)$ and a constant C_3, such that if $\ell \ge \ell_0(\omega)$ and $n = n(\ell)$, then for all $D \in \mathcal{D}_\ell$, we have*

$$|M_n(D) - \mathbb{E}M_n(D)| < \frac{1}{2}\mathbb{E}M_n(D)\,, \tag{2.5}$$

and

$$\sum_{m=\ell}^{n(\ell)} \varphi_m \sum_{J \in \mathcal{D}_m,\, J \subset D} \frac{M_n(J)^2}{(2^{n-\ell} p_n)^2} \le C_3 \varphi_\ell\,. \tag{2.6}$$

Proof: For $m \le n$ and $J \in \mathcal{D}_m$, denote

$$\Delta_n(J) := M_n(J) - \mathbb{E}M_n(J)\,.$$

Also, for $\ell \le n$ and $D \in \mathcal{D}_\ell$, set

$$\Upsilon_n(D) := \sum_{m=\ell}^{n} \varphi_m \sum_{J\in\mathcal{D}_m,\, J\subset D} \Delta_n(J)^2 \,.$$

For $J \in \mathcal{D}_m$, the assumption (2.1) gives $\mathbb{E}[\Delta_n(J)^2] \le \zeta(n) p_n 2^{n-m}$. Therefore,

$$\forall D \in \mathcal{D}_\ell, \quad \mathbb{E}\Upsilon_n(D) = \sum_{m=\ell}^{n} \varphi_m 2^{m-\ell}\zeta(n)p_n 2^{n-m} = 2^{n-\ell}\zeta(n)p_n \sum_{m=\ell}^{n}\varphi_m \,.$$

By (2.3), we then have

$$\mathbb{E}\Upsilon_n(D) \le C2^{n-\ell}p_n^2\varphi_n\zeta(n)/p_n \,.$$

Thus, by (2.4), since $n = n(\ell)$,

$$\mathbb{E}\sum_{D\in\mathcal{D}_\ell} \frac{\Upsilon_n(D)}{(2^{n-\ell}p_n)^2} \le C2^{-\ell} \,.$$

Since the right-hand side is summable in ℓ, we conclude that the summands inside the last expectation tend to 0 a.s. as $\ell \to \infty$. In particular, there exists $\ell_0(\omega) < \infty$ such that for all $\ell \ge \ell_0(\omega)$ and $D \in \mathcal{D}_\ell$, we have

$$\Upsilon_n(D) \le \left(2^{n-\ell}p_n\right)^2 = [\mathbb{E}M_n(D)]^2 \,. \tag{2.7}$$

To deduce (2.5), observe that

$$\Delta_n(D)^2 \le \varphi_\ell^{-1}\Upsilon_n(D) \le \varphi_\ell^{-1}[\mathbb{E}M_n(D)]^2 < \frac{1}{4}[\mathbb{E}M_n(D)]^2 \,.$$

Next, we calculate

$$\sum_{J\in\mathcal{D}_m,\, J\subset D} \frac{[\mathbb{E}M_n(J)]^2}{(2^{n-\ell}p_n)^2} = \sum_{J\in\mathcal{D}_m,\, J\subset D} 2^{2(\ell-m)} = 2^{\ell-m} \,.$$

Therefore, by (2.3),

$$\sum_{m=\ell}^{n}\varphi_m \sum_{J\in\mathcal{D}_m,\, J\subset D} \frac{[\mathbb{E}M_n(J)]^2}{(2^{n-\ell}p_n)^2} = 2^\ell \sum_{m=\ell}^{n} 2^{-m}\varphi_m \le C\varphi_\ell \,. \tag{2.8}$$

Rewrite (2.7) in the form

$$\sum_{m=\ell}^{n} \varphi_m \sum_{J\in\mathcal{D}_m,\, J\subset D} \frac{\Delta_n(J)^2}{(2^{n-\ell}p_n)^2} = \frac{1}{(2^{n-\ell}p_n)^2}\sum_{m=\ell}^{n}\varphi_m \sum_{J\in\mathcal{D}_m,\, J\subset D} \Delta_n(J)^2 \le 1\,. \tag{2.9}$$

Since

$$M_n(J)^2 = [\mathbb{E}M_n(J) + \Delta_n(J)]^2 \le 2[\mathbb{E}M_n(J)]^2 + 2\Delta_n(J)^2\,,$$

adding the inequalities (2.8) and (2.9) yields (2.6), for some constant C_3. □

Proof of Theorem 2.1: We use freely the terminology introduced in the statement of Lemma 2.3. With $\ell_0 = \ell_0(\omega)$ as in the lemma, define inductively $\ell_{k+1} := n(\ell_k)$ for $k \ge 0$. For $D \in \mathcal{D}_{\ell_{k-1}}$ with $k \ge 1$, write

$$Q_k := \mathbb{E}M_{\ell_k}(D) = 2^{\ell_k - \ell_{k-1}} p_{\ell_k}\,,$$

and note that by (2.5),

$$\forall k \ge 1\ \forall D \in \mathcal{D}_{\ell_{k-1}}\,, \quad \frac{1}{2}Q_k \le M_{\ell_k}(D) \le 2Q_k\,. \tag{2.10}$$

Summing this over $D \in \mathcal{D}_{\ell_{k-1}}$ gives

$$\forall k \ge 1 \quad M_{\ell_k}([0,1]) \le 2^{\ell_{k-1}+1} Q_k\,. \tag{2.11}$$

To establish the theorem, we will construct a (random) probability measure ν, supported on $\cap_{k\ge1} A(\ell_k) \subset A$, such that $\mathcal{E}_\varphi(\nu) < \infty$ a.s. To specify ν, it suffices to define $\nu(J)$ consistently for all binary intervals J. Start by assigning the leftmost interval in $\mathcal{D}_{\ell_0}$ full measure, i.e., set $\nu[0, 2^{-\ell_0}] := 1$. Continue inductively:
If $J \in \mathcal{D}_m$ with $\ell_{k-1} < m \le \ell_k$, and $J \subset D$ with $D \in \mathcal{D}_{\ell_{k-1}}$, define

$$\nu(J) := \frac{M_{\ell_k}(J)\nu(D)}{M_{\ell_k}(D)}\,. \tag{2.12}$$

It is straightforward to verify that this assignment is consistent and that ν is supported on $\cap_{k\geq 1} A(\ell_k)$. For $k \geq 2$ and J as in (2.12), two applications of (2.10) and the bound

$$\nu(D) \leq \frac{Z_D}{\min_{\tilde{D}\in\mathcal{D}_{\ell_{k-2}}} M_{\ell_{k-1}}(\tilde{D})},$$

give

$$\nu(J) \leq \frac{2M_{\ell_k}(J)\nu(D)}{Q_k} \leq \frac{4M_{\ell_k}(J)Z_D}{Q_k Q_{k-1}}. \tag{2.13}$$

Now we apply Lemma 2.3. For $k \geq 2$ and $D \in \mathcal{D}_{\ell_{k-1}}$,

$$\begin{aligned} &\sum_{m=\ell_{k-1}}^{\ell_k} \varphi_m \sum_{J\in\mathcal{D}_m,\, J\subset D} \nu(J)^2 \\ &\leq \frac{16 Z_D}{Q_{k-1}^2} \sum_{m=\ell_{k-1}}^{\ell_k} \varphi_m \sum_{J\in\mathcal{D}_m,\, J\subset D} \frac{M_{\ell_k}(J)^2}{Q_k^2} \leq \frac{16 C_3 Z_D}{Q_{k-1}^2} \varphi_{\ell_{k-1}}, \end{aligned} \tag{2.14}$$

by the definition of Q_k and (2.6). Summing this over all $D \in \mathcal{D}_{\ell_{k-1}}$, and then using (2.11) with $k-1$ in place of k, we obtain

$$\begin{aligned} \sum_{m=\ell_{k-1}}^{\ell_k} \varphi_m \sum_{J\in\mathcal{D}_m} \nu(J)^2 &\leq \frac{16 C_3 M_{\ell_{k-1}}([0,1])}{Q_{k-1}^2} \varphi_{\ell_{k-1}} \leq \frac{C_4 2^{\ell_{k-2}}}{Q_{k-1}} \varphi_{\ell_{k-1}} \\ &\leq \frac{C_4 2^{2\ell_{k-2}}}{2^{\ell_{k-1}} p_{\ell_{k-1}}} \varphi_{\ell_{k-1}} \leq C_4 2^{-\ell_{k-2}}, \end{aligned} \tag{2.15}$$

where the last step used (2.4) and the fact that $\zeta \geq 1$. As the right-hand side of (2.15) is summable in k, we conclude that

$$\mathcal{E}_\varphi(\nu) = \sum_{m=0}^{\infty} \varphi_m \sum_{J\in\mathcal{D}_m} \nu(J)^2 < \infty \quad \text{a.s.}$$

By Lemma 2.2 , this completes the proof. $\square$

The next corollary will be used to prove Proposition 1.2 in Section 4. For $K \subset [0,1]$, let $\mathcal{N}_m(K)$ denote the number of intervals in $\mathcal{D}_m$ that intersect K.

Denote by

$$\overline{\dim}_{\mathrm{M}}(K) := \limsup_{m\to\infty} \frac{\log_2 \mathcal{N}_m(K)}{m}\,. \tag{2.16}$$

the **upper Minkowski dimension** of K, also known as the upper box dimension.

The only property of packing dimension $\dim_{\mathrm{P}}$ that we need, is its relation to Minkowski dimension: For any Borel set E,

$$\dim_{\mathrm{P}}(E) = \inf_{E\subset\cup_j E_j} \sup_j \overline{\dim}_{\mathrm{M}}(E_j)\,, \tag{2.17}$$

where the infimum is over all countable covers $\{E_j\}$ of E by closed sets. See Tricot [30] or Falconer [8, Prop. 3.8].

Corollary 2.4. **(i)** *If for each $n \geq 1$ the set V_n is relatively open and dense in $[0,1]^d$, then $\dim_{\mathrm{P}}(\cap_n V_n) = d$.*
(ii) *In the setting of Theorem 2.1, the random set $A = \limsup A(n)$ satisfies $\dim_{\mathrm{P}}(A) = 1$ a.s.*

Proof: **(i)** Let $\{E_j\}$ be a countable collection of closed sets that cover $\cap_n V_n$. Then the union $\{E_j\}_{j\geq 1} \cup \{V_n^c\}_{n\geq 1}$ is a countable cover of $[0,1]^d$ consisting of closed sets. By Baire's theorem, at least one of these closed sets must have nonempty interior in $[0,1]^d$; since each V_n is dense, some E_j must have interior. From (2.17), we conclude that $\dim_{\mathrm{P}}(\cap_n V_n) = d$.
(ii) Denote by $A(n)^\circ$ the interior of $A(n)$, so that $V_n := \cup_{k=n}^{\infty} A(k)^\circ$ is certainly open in $[0,1]$. Fixing a dyadic interval I in $[0,1]$ it is easy to check that Theorem 2.1 applies also when $\mathcal{D}_n$ are the dyadic subintervals of I. Hence, a.s., for each dyadic I, the set $A \cap I$, of positive Hausdorff dimension, is uncountable. intersection Since $A \setminus V_n$ is countable, each V_n is a.s. dense in $[0,1]$, so the assertion

follows from (i). □

3. Localization. Throughout this section, c, c' denote positive, finite constants, independent of ϵ, the values of which may change from line to line, using the notation $a \sim b$ if $\lim_{\epsilon\to 0} a/b = 1$.

To derive lower bounds on the Hausdorff dimension of the sets appearing in Theorem 1.1, as well as for proving (1.7), it is crucial to be able to consider the occupation measure of a ball of radius ϵ over a small time interval (of length δ_ϵ which tends to zero with ϵ), rather than over an interval of constant length.

Surprisingly, it turns out that with only a small loss in probability, we can work with rather short time intervals; the following lemma makes this precise.

LEMMA 3.1 (THE LOCALIZATION LEMMA). *Let $\{W_t\}$ be a Brownian motion in $\mathbb{R}^d$, $d \geq 3$. Write $h(r) := r^2|\log r|$, and $\theta^* := \Lambda_d^{-1} = q_d^2/2$. Finally, denote $\delta_\epsilon := \epsilon^2|\log\epsilon|^6$ and $\beta_\epsilon := 1 - 2|\log\epsilon|^{-2}$. Then for some $0 < c < \infty$, we have*

$$\mathbf{p}_\epsilon := \mathbf{P}(\mu_{\delta_\epsilon}^W(B(0,\epsilon\beta_\epsilon)) \geq ah(\epsilon)) \geq c\epsilon^{a\theta^*} .$$

We did not attempt to optimize the powers of $|\log\epsilon|$ appearing in the definitions of δ_ϵ and β_ϵ. Nevertheless, to appreciate the sharpness of this lemma, recall that by [3], c.f. (3.4) below,

$$\mathbf{P}\Big(\mu_\infty^W(B(0,\epsilon)) \geq ah(\epsilon)\Big) \sim c'\epsilon^{a\theta^*} .$$

Proof: Define

$$\mathcal{T} = \mathcal{T}(\epsilon) := \inf\{s \geq 0 : |W_s| = \epsilon|\log\epsilon|^2\} .$$

By Brownian scaling, we deduce the existence of positive constants c_1, c_2 such that

$$\mathbf{P}(\mathcal{T} > \delta_\epsilon) = \mathbf{P}\Big(\sup_{t\in[0,|\log\epsilon|^2]} |W_t| \le 1\Big) \sim c_1 \exp(-c_2|\log\epsilon|^2)\,. \tag{3.1}$$

Therefore,

$$\begin{aligned}\mathbf{p}_\epsilon &\ge \mathbf{P}(\epsilon^{-2}\int_0^{\mathcal{T}} \mathbf{1}_{\{|W_s|<\epsilon\beta_\epsilon\}} ds \ge a|\log\epsilon|; \mathcal{T} \le \delta_\epsilon) \\ &\ge \mathbf{P}\Big(\epsilon^{-2}\int_0^{\mathcal{T}} \mathbf{1}_{\{|W_s|<\epsilon\beta_\epsilon\}} ds \ge a|\log\epsilon|\Big) - \mathbf{P}(\mathcal{T} > \delta_\epsilon) \end{aligned}\tag{3.2}$$

By (3.1) and (3.2), the lemma will be proved once we establish that

$$\mathbf{P}\Big(\epsilon^{-2}\int_0^{\mathcal{T}} \mathbf{1}_{\{|W_s|<\epsilon\beta_\epsilon\}} ds \ge a|\log\epsilon|\Big) \ge c\epsilon^{a\theta^*}\,. \tag{3.3}$$

To see (3.3), denote by τ_{d-2} the hitting time of the unit sphere in $I\!R^{d-2}$ by Brownian motion, and define

$$I = \epsilon^{-2}\int_0^\infty \mathbf{1}_{\{|W_s|<\epsilon\beta_\epsilon\}} ds\,,$$

$$I^{\mathcal{T}} = \epsilon^{-2}\int_0^{\mathcal{T}} \mathbf{1}_{\{|W_s|<\epsilon\beta_\epsilon\}} ds\,.$$

Recall that, using [3] for the first equality,

$$\frac{\mathbf{P}(\int_0^\infty \mathbf{1}_{\{|W_s|<1\}} ds \ge x)}{e^{-x\theta^*}} = \frac{\mathbf{P}(\tau_{d-2} \ge x)}{e^{-x\theta^*}} \to_{x\to\infty} c\,. \tag{3.4}$$

Therefore, using Brownian scaling and (3.4),

$$\begin{aligned}\mathbf{P}(I \ge a|\log\epsilon|) &= \mathbf{P}\Big((\beta_\epsilon\epsilon)^{-2}\int_0^\infty \mathbf{1}_{\{|W_s|<\epsilon\beta_\epsilon\}} ds \ge \beta_\epsilon^{-2} a|\log\epsilon|\Big) \\ &= \mathbf{P}(\tau_{d-2} \ge \beta_\epsilon^{-2} a|\log\epsilon|) \\ &\sim c\exp\Big(-\theta^* a|\log\epsilon|/(1-2|\log\epsilon|^{-2})^2\Big) \sim c\epsilon^{a\theta^*}\,. \end{aligned}\tag{3.5}$$

Let now $\mathcal{T}' := \inf\{t > \mathcal{T} : |W_t| < \epsilon\}$, and define

$$I_{\mathcal{T}'} = \epsilon^{-2}\int_{\mathcal{T}'}^\infty \mathbf{1}_{\{|W_s|<\epsilon\beta_\epsilon\}} ds\,.$$

Then, $I = I^{\mathcal{T}}\mathbf{1}_{\{\mathcal{T}'=\infty\}} + (I^{\mathcal{T}} + I_{\mathcal{T}'})\mathbf{1}_{\{\mathcal{T}'<\infty\}}$ so that

$$\mathbf{P}(I^{\mathcal{T}} \geq z; \mathcal{T}' = \infty) = \mathbf{P}(I \geq z) - \mathbf{P}(I^{\mathcal{T}} + I_{\mathcal{T}'} \geq z; \mathcal{T}' < \infty) \ . \tag{3.6}$$

Let $\tilde{I}$ be an independent copy of I for a Brownian motion whose expectation when starting at v we denote by $\tilde{\mathbb{E}}^v$. Using symmetry and the strong Markov property we have

$$\begin{aligned}
\mathbf{P}(I^{\mathcal{T}} + I_{\mathcal{T}'} \geq z; \mathcal{T}' < \infty) &= \mathbb{E}\left(\tilde{\mathbb{E}}^{W_{\mathcal{T}'}}\left\{\tilde{I} \geq z - I^{\mathcal{T}}\right\}; \mathcal{T}' < \infty\right) \\
&\leq \mathbb{E}\left(\tilde{\mathbb{E}}\left\{\tilde{I} \geq z - I^{\mathcal{T}}\right\}; \mathcal{T}' < \infty\right) \\
&= \tilde{\mathbb{E}}\left\{\mathbb{E}\left(I^{\mathcal{T}} \geq z - \tilde{I}; \mathcal{T}' < \infty\right)\right\} \\
&= \tilde{\mathbb{E}}\left\{\mathbb{E}\left(\mathbb{E}^{W_{\mathcal{T}}}(T_{B(0,\epsilon)} < \infty);\ I^{\mathcal{T}} \geq z - \tilde{I}\right)\right\} \\
&= |\log\epsilon|^{-2(d-2)}\mathbf{P}(I^{\mathcal{T}} + \tilde{I} \geq z) \\
&\leq |\log\epsilon|^{-2(d-2)}\mathbf{P}(I + \tilde{I} \geq z) \qquad (3.7)
\end{aligned}$$

where $T_{B(0,\epsilon)} = \inf\{t \geq 0 : W_t \in B(0,\epsilon)\}$ denotes the first hitting time of $B(0,\epsilon)$.

Let $\tilde{\tau}_{d-2}$ denote an independent copy of τ_{d-2}, and let q_τ denote their common law. Then, for some constant C independent of z, which may change from line to line,

$$\begin{aligned}
\mathbf{P}(\tau_{d-2} + \tilde{\tau}_{d-2} > z) &= \mathbf{P}(\tau_{d-2} > z) + \int_0^z \mathbf{P}(\tilde{\tau}_{d-2} > z - y) q_\tau(dy) \\
&\leq C\Big[\exp(-z\theta^*) + \int_0^z \exp(-(z-y)\theta^*) q_\tau(dy)\Big] \\
&\leq C\exp(-z\theta^*) + C\int_0^z \exp(-z\theta^*) dy \\
&= C(1+z)\exp(-z\theta^*)\,, \qquad (3.8)
\end{aligned}$$

where the third line came from integration by parts. Hence, by the same argument as in (3.5), for some $c > 0$ and any $\epsilon > 0$ small enough,

$$\mathbf{P}(I + \tilde{I} \geq a|\log \epsilon|) \leq ca|\log \epsilon|\epsilon^{a\theta^*}.$$

Since $2(d-2) > 1$, the inequality (3.3) follows from (3.5), (3.6), (3.7) and the above. □

4. Proof of the lower bound and critical case in Theorem 1.1. The following corollary of Theorem 2.1 and the Localization Lemma will yield the desired lower bound.

Recall that $\theta^* = \Lambda_d^{-1} = q_d^2/2$ denotes the first eigenvalue of the Dirichlet half-Laplacian in the unit ball of $I\!R^{d-2}$.

COROLLARY 4.1. *Let $T \in (0, \infty]$ and $a \in (0, 2\Lambda_d)$. Denote $h(\epsilon) = \epsilon^2|\log \epsilon|$, and consider the set of "thick times"*

$$\mathsf{Thick}_{\geq a} = \{0 \leq t < T \,\Big|\, \limsup_{\epsilon \to 0} \frac{\mu_T^W(B(W_t, \epsilon))}{h(\epsilon)} \geq a\}.$$

Let $\gamma = a\theta^*/2 \in (0,1)$ and $\varphi(r) = r^{1-\gamma}|\log_2 r|^{13}$. Then $\mathcal{H}^\varphi(\mathsf{Thick}_{\geq a}) > 0$ a.s.

Derivation of the lower bound in Theorem 1.1: Assuming for the moment the upper bounds on dimension obtained in Section 5, we may infer that $\dim(\mathsf{Thick}_a) = 1 - \gamma$ as follows (cf. the argument in [13, Pg. 185]). The inequality $\dim(\mathsf{Thick}_{\geq(a+1/n)}) \leq 1 - (a + 1/n)\theta^*/2$ of Section 5 implies that $\mathcal{H}^\varphi(\mathsf{Thick}_{\geq(a+1/n)}) = 0$, and since $\mathsf{Thick}_a = \mathsf{Thick}_{\geq a} - \cup_{n=1}^\infty \mathsf{Thick}_{\geq(a+1/n)}$, Corollary 4.1 shows that $\mathcal{H}^\varphi(\mathsf{Thick}_a) > 0$ which in turn implies that $\dim(\mathsf{Thick}_a) \geq 1 - \gamma$. Using once again the upper bound from Section 5 then completes the proof that $\dim(\mathsf{Thick}_a) = 1 - \gamma$. □

Derivation of the critical case in Theorem 1.1: We now show that $\mathsf{Thick}_{4/q_d^2} \neq ?$; perhaps surprisingly, this can be done by a "soft" argument. For $h > 0$ and $a < 4/q_d^2$, consider the set of approximate thick times

$$\mathsf{Thick}(a,h) := \bigcup_{\epsilon \in (0,h)} \left\{ 0 < t < T \,\Big|\, \frac{\mu_T^W(B(W_t,\epsilon))}{\epsilon^2 |\log \epsilon|} > a \right\}.$$

For any $a < 4/q_d^2$ and $h > 0$, it follows from (1.4) and the Markov property of Brownian motion, that $\mathsf{Thick}(a,h)$ is a.s. dense in $[0,T]$, and it is easy to check that $\mathsf{Thick}(a,h)$ is an open set. Thus fixing sequences $a_n \uparrow 4/q_d^2$ and $h_n \downarrow 0$, Baire's category theorem implies that

$$\cap_n \mathsf{Thick}(a_n, h_n) \neq ?\,.$$

Finally, inspection shows that this intersection coincides with $\mathsf{Thick}_{\geq 4/q_d^2}$, which in turn coincides with Thick_{4/q_d^2} by the remark following Theorem 1.3. $\square$

Proof of Corollary 4.1: Since we are proving a lower bound, we may assume that T is finite; by Brownian scaling, it is enough to consider $T = 2$. Take $\epsilon_n = n^3 2^{-n/2}$, $n = 1, 2, \ldots$ and $\beta_{\epsilon_n} = 1 - 2|\log \epsilon_n|^{-2}$ as in the Localization Lemma. With $I = [t, t + 2^{-n}] \in \mathcal{D}_n$, define $\tilde{I} = [t, t + n^{12} 2^{-n}]$, and let

$$Z_I = 1 \quad \text{iff} \quad \int_{\tilde{I}} \mathbf{1}_{\{|W_s - W_t| < \epsilon_n \beta_{\epsilon_n}\}} ds \geq a h(\epsilon_n)\,.$$

By Lévy's uniform modulus of continuity, there exists an a.s. finite random variable $n_0(\omega)$, such that for all $n \geq n_0(\omega)$,

$$\sup\{|W_t - W_{t'}| \,:\, t, t' \in [0,1], |t - t'| \leq 2^{-n}\} \leq 2\sqrt{2^{-n} \log(2^n)}\,.$$

Therefore, for all $n > n_0(\omega)$, if $I \in \mathcal{D}_n$ and $Z_I = 1$, then

$$\int_{\tilde{I}} \mathbf{1}_{\{|W_s - W_{t'}| < \epsilon_n\}} ds \geq a h(\epsilon_n)$$

for every $t' \in I$. The set A defined in Theorem 2.1 satisfies $A \subset \mathsf{Thick}_{\geq a}$ a.s. (we have taken $T = 2$ rather than $T = 1$ to avoid boundary effects here). The Localization Lemma, Lemma 3.1, shows that for $I \in \mathcal{D}_n$, and all n large enough, $p_n = \mathbf{P}(Z_I = 1) \geq 2^{-a\theta^* n/2}$. Thus, Corollary 4.1 will be established once we verify the variance condition (2.1). For intervals $I, J \in \mathcal{D}_n$ the variables Z_I and Z_J always satisfy $\mathrm{Cov}(Z_I, Z_J) \leq \mathbb{E}(Z_I) = p_n$, and if $\mathrm{dist}(I, J) > n^{12}2^{-n}$, then Z_I and Z_J are independent. Therefore, fixing $m < n$ and $D \in \mathcal{D}_m$, each $I \in \mathcal{D}_n$ satisfies $\mathrm{Cov}(Z_I, M_n(D)) \leq n^{12}p_n$. Consequently

$$\mathrm{Var}(M_n(D)) = \sum_{I \in \mathcal{D}_n, I \subset D} \mathrm{Cov}(Z_I, M_n(D)) \leq 2^{n-m} n^{12} p_n \,.$$

Hence, Theorem 2.1 may be applied (with $p_n \geq 2^{-\gamma n}$ and $\zeta(n) = n^{12}$) to yield the conclusion. $\Box$

Proof of Proposition 1.2: In the course of the proof of Corollary 4.1, we showed that for $a \in (0, 2\Lambda_d)$, the set $\mathsf{Thick}_{\geq a}$ contains a set of the form $A = \limsup A(n)$ that satisfies the hypotheses of Theorem 2.1. Thus, the assertion $\dim_{\mathrm{P}}(\mathsf{Thick}_{\geq a}) = 1$ follows immediately from Corollary 2.4(ii). To handle the critical case $a = 4/q_d^2 = 2\Lambda_d$, observe that in the analysis of that case earlier in this section, we expressed $\mathsf{Thick}_{\geq a}$ as a countable intersection of dense open sets, so $\dim_{\mathrm{P}}(\mathsf{Thick}_{\geq a}) = 1$ by Corollary 2.4(i). (This method also applies to smaller a). Finally, we may deduce (1.6) from the uniform doubling of packing dimension by spatial Brownian motion, established by Perkins-Taylor [18, Cor. 5.8]. $\Box$

5. The upper bound in Theorem 1.1.

In this section we establish the upper bound for (1.3), thus completing the proof of Theorem 1.1. Let $\{W_t\}_{t \geq 0}$ be a Brownian motion in $\mathbb{R}^d$, $d \geq 3$, and $h(\epsilon) = \epsilon^2 |\log \epsilon|$. Set

$$z_T(x, \epsilon) := \mu_T^W(B(x, \epsilon))/h(\epsilon),$$

with $z(x,\epsilon) = z_\infty(x,\epsilon)$. In this section we show that

$$\dim\{x \in B(0,k) \,\Big|\, \limsup_{\epsilon\to 0} z(x,\epsilon) \geq a\} \leq 2 - a\Lambda_d^{-1} \tag{5.1}$$

a.s. for all $a \leq 2\Lambda_d$, $k \in [1,\infty)$. Using $z(x,\epsilon) \geq z_T(x,\epsilon)$, and considering the countable union over $k = 1, 2, \ldots$, will then complete the proof of the upper bound on the dimension of sets in (1.3).

Fix $k \in [1,\infty)$ and $\delta \in (0, 1/5)$. Choose a sequence $\epsilon_n \downarrow 0$ as $n \to \infty$ in such a way that $\epsilon_1 < e^{-2}$ and

$$h(\epsilon_{n+1}) = (1-\delta)h(\epsilon_n). \tag{5.2}$$

For any $a > 0$ let

$$D_a := \{x \in B(0,k) \,\Big|\, \limsup_{n\to\infty} z(x,\epsilon_n) \geq (1-\delta)a\}.$$

Since, for $\epsilon_{n+1} \leq \epsilon \leq \epsilon_n$ we have

$$z(x,\epsilon_n) = \frac{h(\epsilon_{n+1})}{h(\epsilon_n)} \frac{\mu_\infty^W(B(x,\epsilon_n))}{h(\epsilon_{n+1})} \geq (1-\delta) z(x,\epsilon) \tag{5.3}$$

it is easy to see that

$$\{x \in B(0,k) \,\Big|\, \limsup_{\epsilon\to 0} z(x,\epsilon) \geq a\} \subseteq D_a\ .$$

Let $\{x_j : j = 0, 1, \ldots, K_n\}$, with $x_0 = 0$, denote a maximal collection of points in $B(0,k)$ such that $\inf_{\ell\neq j} |x_\ell - x_j| \geq \delta\epsilon_n$. Let $\mathcal{A}_n$ be the set of j, $0 \leq j \leq K_n$, such that

$$\mu_\infty^W(B(x_j, (1+\delta)\epsilon_n)) \geq (1-2\delta) a h(\epsilon_n).$$

We will shortly prove that for any $a > 0$,

$$\mathbb{E}|\mathcal{A}_n| \leq c' \epsilon_n^{(1-4\delta)a\theta^* - 2}. \tag{5.4}$$

Assuming this for the moment, fix $a \leq 2/\theta^*$ and let $\mathcal{V}_{n,j} = B(x_j, \delta\epsilon_n)$. For any $x \in B(0,k)$ there exists $j \in \{0,\ldots,K_n\}$ such that $x \in \mathcal{V}_{n,j}$ and $B(x,\epsilon_n) \subseteq B(x_j,(1+\delta)\epsilon_n)$. Consequently, $\cup_{n\geq m}\cup_{j\in\mathcal{A}_n}\mathcal{V}_{n,j}$ forms a cover of D_a by sets of maximal diameter $2\delta\epsilon_m$. Since $\mathcal{V}_{n,j}$ have diameter $2\delta\epsilon_n$, it follows from (5.4) that for $\gamma = 2-(1-5\delta)a\theta^* > 0$,

$$\mathbb{E}\sum_{n=m}^{\infty}\sum_{j\in\mathcal{A}_n}|\mathcal{V}_{n,j}|^{\gamma} \leq c'(2\delta)^{\gamma}\sum_{n=m}^{\infty}\epsilon_n^{\delta a\theta^*} < \infty\ .$$

Thus, $\sum_{n=m}^{\infty}\sum_{j\in\mathcal{A}_n}|\mathcal{V}_{n,j}|^{\gamma}$ is finite a.s. implying that $\dim(D_a) \leq \gamma$ a.s. Taking $\delta \downarrow 0$ completes the proof of the upper bound (5.1), subject only to (5.4) which we now prove.

Let $\sigma_j = \inf\{t \geq 0 : W_t \in B(x_j,(1+\delta)\epsilon_n)\}$. By the strong Markov property, and [3] (c.f. (3.4)), for some $c = c(\delta,a,d) < \infty$ and all n

$$\begin{aligned}
&\mathbf{P}(\mu_\infty^W(B(x_j,(1+\delta)\epsilon_n)) \geq (1-2\delta)ah(\epsilon_n))\\
&= \mathbf{P}(\mathbb{E}^{W_{\sigma_j}-x_j}\left(\mu_\infty^W(B(0,(1+\delta)\epsilon_n)) \geq (1-2\delta)ah(\epsilon_n)\right)\ ;\sigma_j<\infty)\\
&\leq \mathbf{P}(\mathbb{E}\left(\mu_\infty^W(B(0,(1+\delta)\epsilon_n)) \geq (1-2\delta)ah(\epsilon_n)\right)\ ;\sigma_j<\infty)\\
&\leq c\epsilon_n^{(1-4\delta)a\theta^*}\mathbf{P}(\sigma_j<\infty)
\end{aligned}$$

where the first inequality is due to symmetry. Recall that

$$\mathbf{P}(\sigma_j < \infty) = (\frac{(1+\delta)\epsilon_n}{|x_j|})^{d-2}\wedge 1.$$

Hence, for some $c_1 = c_1(\delta,a,d)$, $c' = c'(\delta,a,d,k) < \infty$ and every n,

$$\begin{aligned}
\text{(5.5)}\qquad \mathbb{E}|\mathcal{A}_n| &= \sum_{j=0}^{K_n}\mathbf{P}(\mu_\infty^W(B(x_j,(1+\delta)\epsilon_n)) \geq (1-2\delta)ah(\epsilon_n))\\
&\leq c_1\epsilon_n^{(1-4\delta)a\theta^*-2}(1+\int_{|x|\leq k}\frac{1}{|x|^{d-2}}\,dx) \leq c'\epsilon_n^{(1-4\delta)a\theta^*-2}
\end{aligned}$$

which completes the proof of (5.4) and consequently that of Theorem 1.1. $\square$

6. Proof of Theorem 1.3. We begin by proving (1.7). To this end, fix $T \in (0,\infty)$, $\delta \in (0,1/4)$ and $a < 2\Lambda_d = 2/\theta^*$ such that $\eta = 2-(1+\delta)a\theta^* > 0$. Choose a sequence $\epsilon_n \downarrow 0$ as in (5.2), noting that for $\epsilon_n \le \epsilon \le \epsilon_{n-1}$ and any $x \in I\!\!R^d$,

$$(1-\delta)z_T(x,\epsilon_n) \le z_T(x,\epsilon) \le (1-\delta)^{-1}z_T(x,\epsilon_{n-1})\,. \tag{6.1}$$

Let $\delta_\epsilon = \epsilon^2|\log\epsilon|^6$, $N_n = [T/\delta_{\epsilon_n}]$, and $t_{i,n} = i\delta_{\epsilon_n}$ for $i = 0,\ldots,N_n-1$. Writing $W_s^t = W_{s+t} - W_t$ it follows that

$$\inf_{\epsilon\in[\epsilon_n,\epsilon_{n-1}]}\sup_{t\in[0,T]} z_T(W_t,\epsilon) \ge (1-\delta)\max_{i=0}^{N_n-1} Z_i^{(n)}\ ,$$

where $Z_i^{(n)} = \mu_{\delta_{\epsilon_n}}^{W^{t_{i,n}}}(B(0,\epsilon_n))/h(\epsilon_n)$ are i.i.d. and by Lemma 3.1, for some $c = c(T) > 0$ and all n large enough,

$$\mathbf{P}(\max_{i=0}^{N_n-1} Z_i^{(n)} \le a) \le (1-\mathbf{p}_{\epsilon_n})^{N_n} \le e^{-c\epsilon_n^{-\eta}}\ .$$

Since ϵ_n^η is summable, applying Borel-Cantelli, then taking $\delta \downarrow 0$ and $a \uparrow 2\Lambda_d$, we see that a.s.

$$\liminf_{\epsilon\to 0}\sup_{t\in[0,T]} z_T(W_t,\epsilon) \ge 2\Lambda_d = 4q_d^{-2}$$

With $S_k(\omega) = \inf\{t : |W_t| \ge k\} \wedge T \in (0,\infty)$ a.s. and $T \mapsto z_T(x,r)$ monotone non-decreasing, it follows that a.s.

$$\liminf_{\epsilon\to 0}\sup_{|x|\le k} z_T(x,\epsilon) \ge \liminf_{\epsilon\to 0}\sup_{t\in[0,S_k(\omega)]} z_{S_k(\omega)}(W_t,\epsilon) \ge 4q_d^{-2}\ .$$

Turning to the proof of the corresponding upper bound, fix $k \in (0,\infty)$, $\delta \in (0,1/5)$ and let $a = (2+\delta)/((1-4\delta)\theta^*) > 2/\theta^*$. Considering the sequence ϵ_n of (5.2) and the sets $\mathcal{A}_n$ as in Section 5, it follows from (5.4) that

$$\sum_{n=1}^\infty \mathbf{P}(|\mathcal{A}_n| \ge 1) \le \sum_{n=1}^\infty I\!\!E|\mathcal{A}_n| \le c'\sum_{n=1}^\infty \epsilon_n^\delta < \infty$$

By Borel-Cantelli, it thus follows that a.s. $\mathcal{A}_n$ is empty for all $n \geq n_0(\omega)$. By the construction of Section 5 the latter event implies that

$$\limsup_{\epsilon\to 0} \sup_{|x|\leq k} z_\infty(x,\epsilon) \leq a \ .$$

Taking $\delta \downarrow 0$ for which $a \downarrow 2/\theta^* = 4q_d^{-2}$, we conclude that a.s.

$$\limsup_{\epsilon\to 0} \sup_{|x|\leq k} z_\infty(x,\epsilon) \leq 4q_d^{-2} \ ,$$

as needed to complete the proof of (1.7).

The left side of (1.8) is monotone in T and by Brownian scaling its law depends only on T/k. Therefore, it suffices to consider $k=1$ and the extreme values $T=1$ and $T=\infty$. Fix $\delta > 0$ and $\epsilon_n = (1-\delta)^n$. (Note, this is different from the ϵ_n used above!). Using the notation

$$\hat{z}_T(x,\epsilon) := \frac{\mu_T^W(B(x,\epsilon))}{(\epsilon^2/|\log\epsilon|)} \ ,$$

it follows that for any $\epsilon \in [\epsilon_n, \epsilon_{n-1}]$ and $x \in I\!\!R^d$

$$\frac{n-1}{n}(1-\delta)^2 \hat{z}_T(x,\epsilon_n) \leq \hat{z}_T(x,\epsilon) \leq \frac{n}{n-1}(1-\delta)^{-2}\hat{z}_T(x,\epsilon_{n-1}) \ .$$

Thus, it suffices for (1.8) to show that for any fixed $\delta \in (0,1/5)$ both the lower bound

$$\liminf_{n\to\infty} \inf_{t\in[0,1]} \hat{z}_1(W_t,\epsilon_n) \geq (1-\delta)^5 \tag{6.2}$$

and the upper bound

$$\limsup_{n\to\infty} \inf_{t\in[0,1]} \hat{z}_\infty(W_t,\epsilon_n) \leq (1+\delta)^5 \ , \tag{6.3}$$

hold a.s.

Our first task in proving (6.2) is to get a good upper bound on the probability of small occupation measure. If $\mu_{[-a,b]}^{\bar{W}}(B(0,\epsilon_n))$ denotes the occupation measure of a two-sided $I\!\!R^d$-valued Brownian motion $\bar{W}$ in $B(0,\epsilon_n)$ during the time interval

$[-a,b]$ with $a,b \geq 0$, then $\mu^{\bar{W}}_{[-a,b]}(B(0,\epsilon_n)) \leq \gamma$ implies that $\bar{\tau}_d(\epsilon_n) \wedge a + \tau_d(\epsilon_n) \wedge b \leq \gamma$, where

$$\bar{\tau}_d(\epsilon) = \inf\{t \geq 0 : |\bar{W}_{-t}| \geq \epsilon\}\ , \qquad \tau_d(\epsilon) = \inf\{t \geq 0 : |\bar{W}_t| \geq \epsilon\}.$$

Taking $\gamma = (1-\delta)^2\epsilon_n^2/|\log\epsilon_n|$, then $a \wedge b \geq (1-\delta)^2\epsilon_n^2/|\log\epsilon_n|$ together with Brownian scaling shows that

$$\begin{aligned}(6.4)\qquad &\mathbf{P}\left(\mu^{\bar{W}}_{[-a,b]}(B(0,\epsilon_n)) \leq (1-\delta)^2\epsilon_n^2/|\log\epsilon_n|\right)\\ &\qquad\qquad \leq \mathbf{P}\left(\bar{\tau}_d(1) + \tau_d(1) \leq (1-\delta)^2/|\log\epsilon_n|\right).\end{aligned}$$

Since $\mathbf{P}(\tau_d(1) \leq x) = \mathbf{P}(\sup_{0\leq t\leq x}|W_t| \geq 1)$, it is well known, see [29, Lemma 6.4], that for $0 < x < 1$

$$c_1 x^{1-d/2} e^{-.5/x} \leq \mathbf{P}(\tau_d(1) \leq x) \leq c_2 x^{1-d/2} e^{-.5/x}. \tag{6.5}$$

This estimate leads, as in the proof of [29, Lemma 6.5], to

$$\mathbf{P}\left(\bar{\tau}_d(1) + \tau_d(1) \leq x\right) \leq e^{-2(1-\delta)/x} \tag{6.6}$$

for any $\delta > 0$ and $x \leq x(\delta)$. Hence, whenever $a \wedge b \geq (1-\delta)^2\epsilon_n^2/|\log\epsilon_n|$,

$$\mathbf{P}\left(\mu^{\bar{W}}_{[-a,b]}(B(0,\epsilon_n)) \leq (1-\delta)^2\epsilon_n^2/|\log\epsilon_n|\right) \leq \epsilon_n^{2/(1-\delta)} \tag{6.7}$$

for all $n \geq n_0(\delta)$, which is the good upper bound we need. In particular, using $\bar{W}^t_s = \bar{W}_{t+s} - \bar{W}_t$ for the time-shifted path, this shows that for all $n \geq n_0(\delta)$,

$$\begin{aligned}(6.8)\qquad &\mathbf{P}(\hat{z}_1(W_t,\epsilon_n) \leq (1-\delta)^2)\\ &= \mathbf{P}\left(\mu^{\bar{W}^t}_{[-t,1-t]}(B(0,\epsilon_n)) \leq (1-\delta)^2\epsilon_n^2/|\log\epsilon_n|\right) \leq \epsilon_n^{2/(1-\delta)}\end{aligned}$$

provided that

$$(1-\delta)^2\epsilon_n^2/|\log\epsilon_n| \leq t \leq 1 - (1-\delta)^2\epsilon_n^2/|\log\epsilon_n| \tag{6.9}$$

On the other hand, if $0 \le t \le 1$ but condition (6.9) does not hold, (i.e. for t close to 0 or 1), we can no longer use the good upper bound (6.8), but must work with the following bound which comes from (6.5):

$$\mathbf{P}(\hat{z}_1(W_t,\epsilon_n) \le (1-\delta)^2) \le \mathbf{P}\left(\tau_d(1) \le (1-\delta)^2/|\log\epsilon_n|\right) \le \epsilon_n^{.5/(1-\delta)^2} \tag{6.10}$$

for all $n \ge n_1(\delta)$, some $n_1(\delta) < \infty$.

To apply these estimates for proving (6.2) take $k = k(\delta) = 20(1-\delta)^2/\delta^2$ to be an integer, $\rho_n = (1-\delta)^2\epsilon_n^2/(|\log\epsilon_n|k(\delta)) = \delta^2\epsilon_{n-1}^2/(20|\log\epsilon_n|)$, $N_n = \lfloor \rho_n^{-1} \rfloor$ and $t_{i,n} = i\rho_n$, $i = 1,\ldots,N_n$. On the one hand, by Lévy's uniform modulus of continuity, we have that a.s. for some finite $n_0 = n_0(\omega) \ge \delta^{-1}$ and all $n \ge n_0$,

$$\max_{i=1}^{N_n} \sup_{|s|<\rho_n} |W_{t_{i,n}+s} - W_{t_{i,n}}| < \delta\epsilon_{n-1} \ ,$$

which implies that

$$\inf_{t\in[0,1]} \hat{z}_1(W_t,\epsilon_{n-1}) \ge (1-\delta)^3 \min_{i=1}^{N_n} \hat{z}_1(W_{t_{i,n}},\epsilon_n) \ . \tag{6.11}$$

On the other hand, we see that condition (6.9) is satisfied by all but the first and last k points of the form $t = t_{i,n}$, $i = 1,\ldots,N_n$. Hence, using the good upper bound (6.8) for those $t_{i,n}$, and the bound (6.10) for the remaining $2k$ $t_{i,n}$'s we have

$$\begin{aligned}\mathbf{P}(\min_{i=1}^{N_n} \hat{z}_1(W_{t_{i,n}},\epsilon_n) \le (1-\delta)^2) &\le \sum_{i=1}^{N_n} \mathbf{P}(\hat{z}_1(W_{t_{i,n}},\epsilon_n) \le (1-\delta)^2)\\ &\le 2k\epsilon_n^{.5/(1-\delta)^2} + N_n\epsilon_n^{2/(1-\delta)} \le \epsilon_n^{2\delta}.\end{aligned} \tag{6.12}$$

Since $\epsilon_n^{2\delta}$ is summable, combining (6.11) and (6.12) yields (6.2) by an application of the Borel-Cantelli Lemma.

Turning to prove (6.3), let now $\gamma_n = (1+\delta)^5\epsilon_n^2/(2|\log\epsilon_n|)$, $\rho_n = \epsilon_n^{2-5.6\delta}$ and n large enough for $\rho_n \ge \gamma_n$. (Our choice of the constant 5.6 will become clear at the end of the proof). Let $\epsilon'_n = (1+\delta)\epsilon_n$ and consider the event $\mathcal{A} = \mathcal{A}^+ \cap \mathcal{A}^-$, where

$$\mathcal{A}^+ = \{\tau_d(\epsilon'_n) \le \gamma_n\,, \inf_{s\in[0,\rho_n]} |W_{\tau_d(\epsilon'_n)+s}| \ge \epsilon_n\,, \,|W_{\tau_d(\epsilon'_n)+\rho_n}| \ge \epsilon_n^{1-\delta}\}$$

and

$$\mathcal{A}^- = \{\bar{\tau}_d(\epsilon'_n) \le \gamma_n\,, \inf_{s\in[0,\rho_n]} |\bar{W}_{-\bar{\tau}_d(\epsilon'_n)-s}| \ge \epsilon_n\,, \,|\bar{W}_{-\bar{\tau}_d(\epsilon'_n)-\rho_n}| \ge \epsilon_n^{1-\delta}\}.$$

By the strong Markov property and symmetry,

$$\begin{aligned}\mathbf{P}(\mathcal{A}^+) &= \mathbf{P}\left(\mathbf{P}^{W_{\tau_d(\epsilon'_n)}}\left(|W_{\rho_n}| \ge \epsilon_n^{1-\delta},\ \inf_{s\in[0,\rho_n]} |W_s| \ge \epsilon_n\right)\,;\tau_d(\epsilon'_n) \le \gamma_n\right)\\ &= \mathbf{P}(\tau_d(\epsilon'_n) \le \gamma_n)\mathbf{P}^{x_0}\left(|W_{\rho_n}| \ge \epsilon_n^{1-\delta},\ \inf_{s\in[0,\rho_n]} |W_s| \ge \epsilon_n\right),\end{aligned}$$

for any x_0 with $|x_0| = \epsilon'_n$.

By Brownian scaling, $\mathbf{P}(\tau_d(\epsilon'_n) \le \gamma_n) = \mathbf{P}(\tau_d(1) \le (1+\delta)^3/(2|\log \epsilon_n|)$, so that using (6.5) and $(1+\delta)^{-3} = 1 - 3\delta + O(\delta^2)$ we get

$$c_3\epsilon_n^{1-2.9\delta} \le \mathbf{P}(\tau_d(\epsilon'_n) \le \gamma_n) \le c_4\epsilon_n^{1-3.1\delta}$$

for some $c_3, c_4 > 0$, δ small and all n large enough. Since

$$\mathbf{P}^x\left(\inf_{s\le 0} |W_s| < \epsilon\right) = \left(\frac{\epsilon}{|x|}\right)^{d-2}, \tag{6.13}$$

whenever $|x| > \epsilon$, we have, with $|x_0| = \epsilon'_n = (1+\delta)\epsilon_n$,

$$\mathbf{P}^{x_0}\left(\inf_{s\le 0} |W_s| \ge \epsilon_n\right) = 1 - (1+\delta)^{-(d-2)} \tag{6.14}$$

hence

$$1 - (1+\delta)^{-(d-2)} \le \mathbf{P}^{x_0}\left(\inf_{s\in[0,\rho_n]} |W_s| \ge \epsilon_n\right) \le 1, \tag{6.15}$$

while

$$\begin{aligned}\mathbf{P}^{x_0}\left(|W_{\rho_n}| \le \epsilon_n^{1-\delta}\right) &= \mathbf{P}\left(|x_0 + \epsilon_n^{1-2.8\delta} W_1| \le \epsilon_n^{1-\delta}\right)\\ &= \mathbf{P}\left(|\epsilon_n^{2.8\delta}(x_0/\epsilon_n) + W_1| \le \epsilon_n^{1.8\delta}\right) \to 0\end{aligned}$$

since $|x_0/\epsilon_n| = 1+\delta$, independent of n. Putting this all together and noting that $\mathbf{P}(\mathcal{A}) = \mathbf{P}(\mathcal{A}^+)\mathbf{P}(\mathcal{A}^-) = \mathbf{P}(\mathcal{A}^+)^2$ shows that

$$c\epsilon_n^{2-5.8\delta} \le \mathbf{P}(\mathcal{A}) \le c'\epsilon_n^{2-6.2\delta} \tag{6.16}$$

for $c, c' > 0$ independent of n.

With $t_{i,n} = 4i\rho_n$ and $N_n = [(4\rho_n)^{-1}] = [0.25\epsilon_n^{-2+5.6\delta}]$, set $\mathcal{A}_i = \mathcal{A} \circ \theta_{t_{i,n}}$, that is, the event $\mathcal{A}$ for the shifted path $W^{t_{i,n}}$ ($\bar{W}^{t_{i,n}}$). By the strong Markov property, for any $i = 1, \ldots, N_n$,

$$\mathbf{P}(\hat{z}_\infty(W_{t_{i,n}}, \epsilon_n) \ge (1+\delta)^5 | \mathcal{A}_i) \le 2 \max_{|x_0| \ge \epsilon_n^{1-\delta}} \mathbf{P}^{x_0}(\inf_{t \ge 0} |W_t| < \epsilon_n) \le 2\epsilon_n^{(d-2)\delta},$$

where (6.13) was used in the second inequality. Hence, by the independence of the events $\{\mathcal{A}_i\}_{i=1}^{N_n}$,

$$\begin{aligned}
&\mathbf{P}(\min_{i=1}^{N_n} \hat{z}_\infty(W_{t_{i,n}}, \epsilon_n) \ge (1+\delta)^5) \\
&\le (1-\mathbf{P}(\mathcal{A}))^{N_n} + \sum_{i=1}^{N_n} \mathbf{P}(\hat{z}_\infty(W_{t_{i,n}}, \epsilon_n) \ge (1+\delta)^5, \mathcal{A}_i) \\
&\le e^{-\mathbf{P}(\mathcal{A})N_n} + \sum_{i=1}^{N_n} \mathbf{P}(\hat{z}_\infty(W_{t_{i,n}}, \epsilon_n) \ge (1+\delta)^5 \,|\, \mathcal{A}_i)\mathbf{P}(\mathcal{A}_i) \\
&\le e^{-c\epsilon_n^{-.2\delta}} + c'\epsilon_n^{-.6\delta}\epsilon_n^{(d-2)\delta} \\
&\le e^{-c\epsilon_n^{-.2\delta}} + c'\epsilon_n^{.4\delta}
\end{aligned}$$

and (6.3) follows by an application of the Borel-Cantelli Lemma. (One can see now the reason for choosing 5.6 above. With more care, we could have chosen $\rho_n = \epsilon_n^{2-q\delta}$ with $5 < q < 6$). This completes the proof of Theorem 1.3. $\square$

7. Proof of Theorem 1.4. For any Borel function $f : [a,b] \to I\!R^d$, we use $\mu_{a,b}^f$ to denote its occupation measure:

$$\mu_{a,b}^f(A) = \int_a^b \mathbf{1}_A(f_t)\,dt$$

for all Borel sets $A \subseteq I\!R^d$. We use the abbreviations $\mu_T^f = \mu_{0,T}^f$ and $\bar{\mu}_T^f = \mu_{-T,T}^f$.

As a first step in proving Theorem 1.4, we rewrite things so that we deal only with occupation measures of $B(0,1)$. Writing $W_t^\epsilon = \epsilon^{-1} W_{\epsilon^2 t}$ and $W_s^{\epsilon,t} = W_{t+s}^\epsilon - W_t^\epsilon$ with similar notation for $\bar{W}$ we have

$$\begin{aligned}\mu_1^W\left(B(W_{\epsilon^2 t},\epsilon)\right) &= \int_0^1 \mathbf{1}_{\{|W_s - W_{\epsilon^2 t}|\le\epsilon\}}\,ds = \epsilon^2\int_0^{1/\epsilon^2}\mathbf{1}_{\{|W_{\epsilon^2 s}-W_{\epsilon^2 t}|\le\epsilon\}}\,ds\\ &= \epsilon^2\int_0^{1/\epsilon^2}\mathbf{1}_{\{|W_s^\epsilon - W_t^\epsilon|\le 1\}}\,ds = \epsilon^2\int_0^{1/\epsilon^2}\mathbf{1}_{B(0,1)}(W_s^\epsilon - W_t^\epsilon)\,ds\end{aligned}$$

and consequently

$$\begin{aligned}\int_0^1 e^{\theta\mu_1^W(B(W_t,\epsilon))/\epsilon^2}\,dt &= \epsilon^2\int_0^{1/\epsilon^2} e^{\theta\mu_1^W\left(B(W_{\epsilon^2 t},\epsilon)\right)/\epsilon^2}\,dt\\ &= \epsilon^2\int_0^{1/\epsilon^2}\exp\left(\theta\int_0^{1/\epsilon^2}\mathbf{1}_{B(0,1)}(W_s^\epsilon - W_t^\epsilon)\,ds\right)dt \qquad (7.1)\\ &\le \epsilon^2\int_0^{1/\epsilon^2}\exp\left(\theta\int_{-\infty}^{\infty}\mathbf{1}_{B(0,1)}(\bar{W}_s^\epsilon - \bar{W}_t^\epsilon)\,ds\right)dt\\ &= \epsilon^2\int_0^{1/\epsilon^2}\exp\left(\theta\int_{-\infty}^{\infty}\mathbf{1}_{B(0,1)}(\bar{W}_s^{\epsilon,t})\,ds\right)dt\\ &= \epsilon^2\int_0^{1/\epsilon^2} e^{\theta\bar{\mu}_\infty^{\bar{W}^{\epsilon,t}}(B(0,1))}\,dt.\end{aligned}$$

Hence for each $\theta < q_d^2/2$ and any subsequence $\epsilon_m \to 0$, in order to show that

$$\limsup_{m\to\infty}\int_0^1 e^{\theta\mu_1^W(B(W_t,\epsilon_m))/\epsilon_m^2}\,dt \le I\!E\left(e^{\theta\bar{\mu}_\infty^W(B(0,1))}\right) \qquad a.s. \tag{7.2}$$

it suffices to show that

$$\lim_{m\to\infty}\epsilon_m^2\int_0^{1/\epsilon_m^2} e^{\theta\bar{\mu}_\infty^{W^{\epsilon_m,t}}(B(0,1))}\,dt = I\!E\left(e^{\theta\bar{\mu}_\infty^W(B(0,1))}\right) \qquad a.s. \tag{7.3}$$

For any $1 < p < 2$ such that $p\theta < q_d^2/2$, (7.3) will follow with $\epsilon_m = m^{-2/(p-1)}$ from the Borel-Cantelli lemma, Chebycheff's inequality and the following lemma. For notational convienience be shall sometimes write $\bar{W}(n^{-1},t)$ for $\bar{W}^{n^{-1},t}$.

Lemma 7.1. *For $\theta < q_d^2/2$, there exists $c = c_{d,\theta}$ finite, such that for all n,*

$$\|\frac{1}{n^2}\int_0^{n^2} e^{\theta\bar{\mu}_\infty^{\bar{W}(n^{-1},t)}(B(0,1))}\,dt - \frac{1}{n^2}\int_0^{n^2} e^{\theta\bar{\mu}_n^{\bar{W}(n^{-1},t)}(B(0,1))}\,dt\|_1 \le cn^{-(d/2-1)}, \tag{7.4}$$

and for any $1 < p < 2$ such that $p\theta < q_d^2/2$, there exists $c = c_{p,d,\theta}$ finite, such that for all n,

$$\|\frac{1}{n^2}\int_0^{n^2} e^{\theta\bar{\mu}_n^{\bar{W}(n^{-1},t)}(B(0,1))}\,dt - \mathbb{E}\left(e^{\theta\bar{\mu}_\infty^{\bar{W}}(B(0,1))}\right)\|_p \le cn^{-(1-1/p)}. \tag{7.5}$$

Before proving this lemma we first use it to show that for any $1 < p < 2$ such that $p\theta < q_d^2/2$, and with $\epsilon_m = m^{-2/(p-1)}$,

$$\liminf_{m\to\infty}\int_0^1 e^{\theta\mu_1^W(B(W_t,\epsilon_m))/\epsilon_m^2}\,dt \ge \mathbb{E}\left(e^{\theta\bar{\mu}_\infty^{\bar{W}}(B(0,1))}\right) \qquad a.s. \tag{7.6}$$

Note that for any $n \le t \le n^2 - n$

$$\int_0^{n^2} \mathbf{1}_{B(0,1)}(W_s^{n^{-1}} - W_t^{n^{-1}})\,ds = \int_{-t}^{n^2-t} \mathbf{1}_{B(0,1)}(\bar{W}_s^{n^{-1},t})\,ds$$
$$\ge \int_{-n}^{n} \mathbf{1}_{B(0,1)}(\bar{W}_s^{n^{-1},t})\,ds = \bar{\mu}_n^{\bar{W}(n^{-1},t)}\left(B(0,1)\right).$$

Hence from (7.1),

$$\int_0^1 e^{\theta\mu_1^W\left(B(W_t,n^{-1})\right)/n^{-2}}\,dt$$
$$= \frac{1}{n^2}\int_0^{n^2}\exp\left(\theta\int_0^{n^2}\mathbf{1}_{B(0,1)}(W_s^{n^{-1}} - W_t^{n^{-1}})\,ds\right)\,dt$$
$$\ge \frac{1}{n^2}\int_n^{n^2-n} e^{\theta\bar{\mu}_n^{\bar{W}(n^{-1},t)}(B(0,1))}\,dt\,.$$

(7.6) then follows by using Lemma 7.1 as before and noting that

$$\|\frac{1}{n^2}\int_0^{n^2} e^{\theta\bar{\mu}_n^{\bar{W}(n^{-1},t)}(B(0,1))}\,dt - \frac{1}{n^2}\int_n^{n^2-n} e^{\theta\bar{\mu}_n^{\bar{W}(n^{-1},t)}(B(0,1))}\,dt\|_1$$
$$\le \|\frac{1}{n^2}\int_0^{n} e^{\theta\bar{\mu}_n^{\bar{W}(n^{-1},t)}(B(0,1))}\,dt\|_1 + \|\frac{1}{n^2}\int_{n^2-n}^{n^2} e^{\theta\bar{\mu}_n^{\bar{W}(n^{-1},t)}(B(0,1))}\,dt\|_1$$

$$\leq \frac{1}{n^2}\int_0^n \|e^{\theta\bar{\mu}_n^{\bar{W}(n^{-1},t)}(B(0,1))}\|_1\,dt + \frac{1}{n^2}\int_{n^2-n}^{n^2} \|e^{\theta\bar{\mu}_n^{\bar{W}(n^{-1},t)}(B(0,1))}\|_1\,dt$$
$$\leq 2n^{-1}\|e^{\theta\bar{\mu}_\infty^{\bar{W}}(B(0,1))}\|_1\,.$$

Since $\mu_1^W(B(W_t,\epsilon))$ is monotone in ϵ and $\lim_{m\to\infty}\epsilon_{m+1}/\epsilon_m = 1$, the proof of Theorem 1.4 now follows from (7.2), (7.6) and a simple interpolation argument.

Proof of Lemma 7.1: (7.4) will follow from

$$\|\frac{1}{n^2}\int_0^{n^2} e^{\theta\bar{\mu}_\infty^{\bar{W}(n^{-1},t)}(B(0,1))}\,dt - \frac{1}{n^2}\int_0^{n^2} e^{\theta\bar{\mu}_n^{\bar{W}(n^{-1},t)}(B(0,1))}\,dt\|_1$$
$$\leq \frac{1}{n^2}\int_0^{n^2} \|e^{\theta\bar{\mu}_\infty^{\bar{W}(n^{-1},t)}(B(0,1))} - e^{\theta\bar{\mu}_n^{\bar{W}(n^{-1},t)}(B(0,1))}\|_1\,dt$$

and the following lemma.

Lemma 7.2. *For any $\theta < q_d^2/2$, there exists $c = c_{d,\theta}$ finite such that for any $\epsilon > 0$,*

$$\|e^{\theta\bar{\mu}_\infty^{\bar{W}}(B(0,1))} - e^{\theta\bar{\mu}_{1/\epsilon}^{\bar{W}}(B(0,1))}\|_1 \leq c\epsilon^{d/2-1}.$$

As for (7.5), we first rewrite

$$\frac{1}{n^2}\int_0^{n^2} e^{\theta\bar{\mu}_n^{\bar{W}(n^{-1},t)}(B(0,1))}\,dt = \frac{1}{n}\sum_{k=0}^{n-1}\frac{1}{n}\int_{kn}^{(k+1)n} e^{\theta\bar{\mu}_n^{\bar{W}(n^{-1},t)}(B(0,1))}\,dt$$
$$= \frac{1}{n}\sum_{k=0}^{n-1} I_{n,k}$$

where

$$I_{n,k} = \frac{1}{n}\int_{kn}^{(k+1)n} e^{\theta\bar{\mu}_n^{\bar{W}(n^{-1},t)}(B(0,1))}\,dt.$$

Unraveling the definitions we see that for each fixed n, the $I_{n,k}\,;\ 0\leq k\leq n$ are identically distributed, and $I_{n,k}$ is measurable with respect to the σ-algebra generated by $\{\bar{W}_{t+s}^{n^{-1}} - \bar{W}_t^{n^{-1}}\,;\ kn\leq t\leq (k+1)n\,;\ -n\leq s\leq n\}$. Hence $I_{n,k}, I_{n,k'}$

are independent as soon as $|k-k'|\geq 3$. Thus we can write

$$\frac{1}{n}\sum_{k=0}^{n} I_{n,k} = \frac{1}{n}\sum_{k=0}^{n/3} I_{n,3k} + \frac{1}{n}\sum_{k=0}^{n/3} I_{n,1+3k} + \frac{1}{n}\sum_{k=0}^{n/3} I_{n,2+3k}$$

where each of the three sums on the right hand side is now a sum of i.i.d. random variables. Furthermore

$$\mathbb{E}\left(I_{n,k}\right) = \frac{1}{n}\int_{kn}^{(k+1)n} \mathbb{E}\left(e^{\theta \bar{\mu}_n^{W(n^{-1},t)}(B(0,1))}\right) dt = \mathbb{E}\left(e^{\theta \bar{\mu}_n^{W}(B(0,1))}\right).$$

Using Lemma 7.2, to complete the proof of Lemma 7.1 it now suffices to note that for any $1<p<2$ such that $p\theta < q_d^2/2$ we have the following bounds, where the first inequality comes from the Marcinkiewicz-Zygmund inequality (see for example [25, Pg. 341], where our condition $p\theta < q_d^2/2$ guarantees that $I_{n,k}\in L^p$), and the second inequality comes from the fact that $|a+b|^{p/2}\leq |a|^{p/2}+|b|^{p/2}$ (since $p<2$):

$$\begin{aligned}
&\mathbb{E}\left(|\frac{1}{n/3}\sum_{k=0}^{n/3}\left(I_{n,i+3k}-\mathbb{E}\left(I_{n,i+3k}\right)\right)|^p\right) \\
&\leq \frac{c}{n^p}\mathbb{E}\left(|\sum_{k=0}^{n/3}\left(I_{n,i+3k}-\mathbb{E}\left(I_{n,i+3k}\right)\right)^2|^{p/2}\right) \leq cn^{-(p-1)}
\end{aligned}$$

for $i=0,1,2$.

Proof of Lemma 7.2: Let $p_r(x) = (2\pi r)^{-d/2}\exp(-|x|^2/2r)$ and $u^0(x) = \int_0^\infty p_r(x)dr = \frac{c_d}{|x|^{d-2}}$ denote the zero-potential of W_t. Let Λ_d denote the norm of

$$Kf(x) = \int_{B(0,1)} u^0(x-y)f(y)\,dy$$

considered as an operator from $L^2\left(B(0,1),\,dx\right)$ to itself, recalling from [3] that $\Lambda_d^{-1} = q_d^2/2$ is the first eigenvalue of $(1/2)\Delta$ in the unit ball of $\mathbb{R}^{d-2}$ (not $\mathbb{R}^d$!) with Dirichlet boundary conditions. We claim that $K^i u^0\in L^2(B(0,1),\,dx)$ for sufficiently large i, ($i=[d/2]$ will do). To see this, note that $u^0\in L^2(B(0,1),\,d^3x)$ for $d=3$, while for $d=4$, by scaling, $Ku^0(x)\leq c\int|x-y|^{-2}|y|^{-2}\,d^4y\leq$

$c\log(1/|x|) \in L^2(B(0,1), d^4x)$. When $d > 4$ we first note, again by scaling, that $\int |x-y|^{-(d-2)}|y|^{-k}\,d^dy \le c|x|^{-(k-2)}$ for any $k > 2$, and we can repeat this argument until we find that $K^j u^0(x) \le c|y|^{-k}$ with $1 \le k \le 2$ for some j. Noting that for such k we have $|y|^{-k} \in L^2(B(0,1), d^dx)$ completes the verification of our claim. It follows from this in particular, for some $\kappa_d < \infty$ and all i

$$\int_{B(0,1)} K^i u^0(x)dx \le \kappa_d(\Lambda_d)^i\,. \tag{7.7}$$

Setting $t_0 = 0$, $x_0 = 0$, we first bound the following moments for $m = 1, 2, \dots$

$$\begin{aligned}
&\frac{1}{m!}\mathbb{E}\left(\int_{1/\epsilon}^{\infty} \mathbf{1}_{B(0,1)}(W_r)\,dr\{\int_0^{\infty} \mathbf{1}_{B(0,1)}(W_s)\,ds\}^{m-1}\right)\\
&= \frac{1}{m}\sum_{i=1}^m \int_{B(0,1)^m}\int_{\substack{0\le t_1\le\cdots\le t_m<\infty\\ \epsilon^{-1}\le t_i}} \prod_{j=1}^m p_{t_j-t_{j-1}}(x_j-x_{j-1})\,dt_1\cdots dt_m\,dx_1\cdots dx_m\\
&= \frac{1}{m}\sum_{i=1}^m \int_{B(0,1)^m}\int_{\epsilon^{-1}\le\sum_{j=1}^i r_j} \prod_{j=1}^m p_{r_j}(x_j-x_{j-1})\,dr_1\cdots dr_m\,dx_1\cdots dx_m\\
&\le \sum_{i=1}^m \int_{B(0,1)^m}\int_{(m\epsilon)^{-1}\le r_i} \prod_{j=1}^m p_{r_j}(x_j-x_{j-1})\,dr_1\cdots dr_m\,dx_1\cdots dx_m\\
&= \sum_{i=1}^m \int_{B(0,1)^m} \prod_{\substack{j=1\\ j\ne i}}^m u^0(x_j-x_{j-1})\int_{(m\epsilon)^{-1}}^{\infty} p_{r_i}(x_i-x_{i-1})\,dr_i\,dx_1\cdots dx_m\\
&\le \left(\int_{(m\epsilon)^{-1}}^{\infty} p_r(0)\,dr\right)\sum_{i=1}^m \int_{B(0,1)^m} \prod_{\substack{j=1\\ j\ne i}}^m u^0(x_j-x_{j-1})\,dx_1\cdots dx_m\\
&= k_d(m\epsilon)^{d/2-1}\Big[(1, K^{m-1}1)_{B(0,1)}\\
&\qquad\qquad + \sum_{i=2}^m \big(\int_{B(0,1)} K^{i-2}u^0(x_{i-1})dx_{i-1}\big)(1, K^{m-i}1)_{B(0,1)}\Big]\\
&\le c_d\epsilon^{d/2-1}m^{d/2}\Lambda_d^{m-1}\,,
\end{aligned}$$

where $(\cdot,\cdot)_{B(0,1)}$ denotes the inner product in $L^2(B(0,1), dx)$ and (7.7) was used in the last inequality. With c_d independent of m and $g_{d,\theta} = \mathbb{E}(e^{\theta\mu_\infty^W(B(0,1))})$

finite, it follows that

$$
\begin{aligned}
&\|e^{\theta\bar{\mu}^{\bar{W}}_{\infty}(B(0,1))} - e^{\theta\bar{\mu}^{\bar{W}}_{1/\epsilon}(B(0,1))}\|_1 \\
&= \mathbb{E}\left(e^{\theta\bar{\mu}^{\bar{W}}_{\infty}(B(0,1))}\right) - \mathbb{E}\left(e^{\theta\bar{\mu}^{\bar{W}}_{1/\epsilon}(B(0,1))}\right) \\
&= \{\mathbb{E}\left(e^{\theta\mu^{W}_{\infty}(B(0,1))}\right)\}^2 - \{\mathbb{E}\left(e^{\theta\mu^{W}_{1/\epsilon}(B(0,1))}\right)\}^2 \\
&\le 2g_{d,\theta}\,|\mathbb{E}\left(e^{\theta\mu^{W}_{\infty}(B(0,1))}\right) - \mathbb{E}\left(e^{\theta\mu^{W}_{1/\epsilon}(B(0,1))}\right)| \\
&\le 2g_{d,\theta}\theta\,\mathbb{E}\left(|\mu^{W}_{\infty}(B(0,1)) - \mu^{W}_{1/\epsilon}(B(0,1))|e^{\theta\mu^{W}_{\infty}(B(0,1))}\right) \\
&= 2g_{d,\theta}\sum_{m=0}^{\infty}\frac{\theta^{m+1}}{m!}\mathbb{E}\left(\int_{1/\epsilon}^{\infty}\mathbf{1}_{B(0,1)}(W_r)\,dr\{\int_0^{\infty}\mathbf{1}_{B(0,1)}(W_s)\,ds\}^m\right) \\
&\le 2\theta g_{d,\theta}c_d\epsilon^{d/2-1}\sum_{m=0}^{\infty}(m+1)^{d/2+1}(\theta\Lambda_d)^m \le c_{d,\theta}\epsilon^{d/2-1}
\end{aligned}
$$

for any $\theta < \Lambda_d^{-1}$, as needed to complete the proof of Lemma 7.2 and hence of Theorem 1.4. □

8. Proof of Corollary 1.5. The lower bound in Corollary 1.5 is an immediate consequence of (1.10) and Chebycheff's inequality. Turning to the corresponding upper bound, fix $a \in (0, 2/\theta^*)$. Choosing $\delta \in (0, 1/4)$ such that $\eta = 2 - a\theta^*(1 + 3\delta) > 0$ and ϵ_n as in (5.2), leads (see (6.1)) to

$$
\begin{aligned}
(8.1)\qquad &\limsup_{\epsilon\to 0}\frac{\log \mathcal{L}eb\left\{0 \le t \le 1 \,\middle|\, z_1(W_t,\epsilon) \ge a\right\}}{\log\epsilon} \\
&\le \limsup_{n\to\infty}\frac{\log \mathcal{L}eb\left\{0 \le t \le 1 \,\middle|\, z_1(W_t,\epsilon_n) \ge \frac{a}{1-\delta}\right\}}{\log\epsilon_{n-1}}\,.
\end{aligned}
$$

Let $W^t_s = W_{t+s} - W_t$, $\delta_n = \epsilon_n^2|\log\epsilon_n|^6$ and $\beta_n = 1 - 2|\log\epsilon_n|^{-2}$. The random variables $Y^{(n)}_i = \mu^{W^{i\delta_n}}_{\delta_n}(B(0,\beta_n\epsilon_n))/h(\epsilon_n)$, $i = 1,\ldots,\delta_n^{-1}-1$ are i.i.d. The Localization Lemma implies that for some $c > 0$ and all n large enough,

$$p^*_n := \mathbf{P}(Y^{(n)} \ge a/(1-\delta)) \ge c\epsilon_n^{a\theta^*(1+2\delta)}\,.$$

Thus, by standard tail estimates for the Binomial$(\delta_n^{-1}-1, p_n^*)$, for all n large enough,

$$\mathbf{P}(|\{i : Y_i^{(n)} \geq a/(1-\delta)\}| \leq \epsilon_n^{-\eta}) \leq \exp(-\epsilon_n^{-\eta}) \ ,$$

since $(1-\delta)^{-1} \leq 1+2\delta$. It follows that a.s., for all $n \geq n_0(\omega, \delta, a)$,

$$|\{i : Y_i^{(n)} \geq a/(1-\delta)\}| \geq \epsilon_n^{-\eta} \ . \tag{8.2}$$

Taking $\rho_n = \epsilon_n^2/|\log \epsilon_n|^6$, by Lévy's uniform modulus of continuity, we have that a.s. for some $n_1 = n_1(\omega) < \infty$ and all $n \geq n_1$,

$$\max_{i=1}^{\delta_n^{-1}-1} \sup_{|s|<\rho_n} |W_{i\delta_n+s} - W_{i\delta_n}| < (1-\beta_n)\epsilon_n \ ,$$

which together with (8.2) implies that a.s. for any $n \geq n_2(\omega, \delta, a)$,

$$\mathcal{L}eb\Big\{0 \leq t \leq 1 \,\Big|\, z_1(W_t, \epsilon_n) \geq a/(1-\delta)\Big\} \geq \rho_n |\{i : Y_i^{(n)} \geq a/(1-\delta)\}| \geq \epsilon_n^{2-\eta+\delta} \ .$$

In view of (8.1), we have a.s.

$$\limsup_{\epsilon \to 0} \frac{\log \mathcal{L}eb\Big\{0 \leq t \leq 1 \,\Big|\, z_1(W_t, \epsilon) \geq a\Big\}}{\log \epsilon} \leq (1-2\delta)^{-1/2}(2-\eta-\delta) \ .$$

To complete the proof consider $\delta \downarrow 0$, for which $2-\eta-\delta \to a\theta^*$. $\square$

9. Large occupation measure at all scales.

Proof of Theorem 1.6: For $k \in (1, \infty)$, $T < \infty$, let $\Gamma_k = \{x : |x| \in [1/k, k]\}$ and

$$D_a := \{x \in \Gamma_k \,\Big|\, \liminf_{\epsilon \to 0} \frac{\mu_T^W(B(x,\epsilon))}{\epsilon^2} \geq a\} \ .$$

(We work with the annulus Γ_k rather than the ball $B(0,k)$ because the basic bound we will use, Lemma 9.2, blows up at the origin).

Fix $\delta > 0$ and let $b = 1+\delta > 1$. Set $\eta_n = 2^{-n}$ and $\delta_n = \eta_n^{1-b^{-1}}$ for $n = 1, 2, \ldots$. Let $\{x_j : j = 1, \ldots, K_n\}$, $K_n \leq c(\delta, k, d)\eta_n^{-d}$, denote a maximal collection of points in Γ_k such that $\inf_{\ell\neq j}|x_\ell - x_j| \geq \delta\eta_n$. Let $\mathcal{H}_n = \mathcal{H}_n(a,\delta,T)$ be the set of j, $1 \leq j \leq K_n$, such that

$$\inf_{\epsilon\in[\eta_n,\delta_n]} \frac{\mu_T^W(B(x_j, b\epsilon))}{\epsilon^2} \geq \frac{a}{b}. \tag{9.1}$$

We will shortly prove that for any $\gamma > 0$ we can find $\delta > 0$ such that for some $c = c(a,\delta,T) < \infty$ and all n,

$$\mathbb{E}|\mathcal{H}_n| \leq c\eta_n^{I_d(a)-2-\gamma} \tag{9.2}$$

where $I_d(v)$ is defined in (1.12). Assuming this for the moment, let $\mathcal{U}_{n,j} = B(x_j, \delta\eta_n)$. Then, for any $x \in \Gamma_k$ there exists $j \in \{1,\ldots,K_n\}$ such that $x \in \mathcal{U}_{n,j}$ and $B(x,\epsilon) \subseteq B(x_j, \epsilon + \delta\eta_n) \subseteq B(x_j, b\epsilon)$ for all $\epsilon \geq \eta_n$. If $x \in D_a$ then a.s. for some $m_1(\omega, x, b) < \infty$ and all $n \geq m_1$,

$$\inf_{\epsilon\in[\eta_n,\delta_n]} \frac{\mu_T^W(B(x,\epsilon))}{\epsilon^2} \geq \frac{a}{b}\ .$$

Therefore, $\cup_{n\geq m}\cup_{j\in\mathcal{H}_n}\mathcal{U}_{n,j}$ forms a $2\delta\eta_m$-cover of D_a for any $m \geq 1$. Since $\mathcal{U}_{n,j}$ has diameter $2\delta\eta_n$, it follows from (9.2) that

$$\begin{aligned} &\mathbb{E}\sum_{n=m}^{\infty}\sum_{j\in\mathcal{H}_n}|\mathcal{U}_{n,j}|^{2-I_d(a)+2\gamma} \\ &= \sum_{n=m}^{\infty}\mathbb{E}|\mathcal{H}_n|(2\delta\eta_n)^{2-I_d(a)+2\gamma} \leq c_2\sum_{n=m}^{\infty}\eta_n^{\gamma} < \infty. \end{aligned} \tag{9.3}$$

Thus, $\sum_{n=m}^{\infty}\sum_{j\in\mathcal{H}_n}|\mathcal{U}_{n,j}|^{2-I_d(a)+2\gamma}$ is finite a.s. implying that $\dim(D_a) \leq 2 - I_d(a) + 2\gamma$ a.s. for any $T < \infty$, $\gamma > 0$. Since a.s. there exists $T_k = T_k(\omega)$ finite, such that $|W_t| \geq (k+1)$ for any $t \geq T_k$, obviously a.s. also

$$\dim(\{x \in \Gamma_k \mid \liminf_{\epsilon\to 0}\frac{\mu_\infty^W(B(x,\epsilon))}{\epsilon^2} \geq a\}) \leq 2 - I_d(a) + 2\gamma\ .$$

Taking $\gamma \downarrow 0$ and considering the countable union over $k = 1, 2\ldots$ completes the proof of (1.14).

To get our upper bound on packing dimension, denote by $D^m(a/b)$ the set of points $x \in I\!R^d$ such that for all $n \geq m$, we have

$$\inf_{\epsilon \in [\eta_n, \delta_n]} \frac{\mu_T^W(B(x,\epsilon))}{\epsilon^2} \geq \frac{a}{b} .$$

Clearly, $\cup_{j \in \mathcal{H}_n} \mathcal{U}_{n,j}$ forms a $2\delta\eta_n$-cover of $D^m(a/b)$ for any $n \geq m$. Thus, from (9.3)

$$\lim_{n\to\infty} \sum_{j \in \mathcal{H}_n} |\mathcal{U}_{n,j}|^{2-I_d(a)+2\gamma} = 0 \quad \text{a.s.} \tag{9.4}$$

Denote by $\mathcal{N}(A,\epsilon)$ the minimal cardinality of a collection of balls of radius ϵ that covers A. Recall that $\overline{\dim}_{\text{M}}(A)$, the **upper Minkowski dimension** of a set A (also known as the *upper box-counting dimension*), may be defined by

$$\overline{\dim}_{\text{M}}(A) = \limsup_{\epsilon\to 0} \frac{\log \mathcal{N}(A,\epsilon)}{|\log \epsilon|} \; ; \tag{9.5}$$

see [8, (3.5)]. From (9.4) we may deduce that $\overline{\dim}_{\text{M}}(D^m(a/b)) \leq 2 - I_d(a) + 2\gamma$. Since $b > 1$, necessarily

$$D_a \subset \cup_{m \geq 1} D^m(a/b) \, , \tag{9.6}$$

and the upper bound $\dim_{\text{P}}(D_a) \leq 2 - I_d(a) + 2\gamma$ a.s. follows by [8, Prop. 3.8]. This completes the proof of (1.15).

We next recall that $I_d(v)$ of (1.12) is strictly increasing in $v \geq 2/(d-2)$, whereas $I_d(C_d) = 2$. Hence, fixing $a > C_d$, we may and shall fix $\gamma > 0$ such that $I_d(a) - 2 - \gamma > 0$. Then, by (9.2), for any $\delta > 0$ sufficiently small

$$\sum_{n=1}^{\infty} \mathbf{P}(|\mathcal{H}_n| \geq 1) \leq \sum_{n=1}^{\infty} I\!E|\mathcal{H}_n| \leq c_1 \sum_{n=1}^{\infty} \eta_n^{I_d(a)-2-\gamma} < \infty.$$

Thus, by Borel-Cantelli, it follows that a.s. $\mathcal{H}_n$ is empty for all $n \geq m_2(\omega)$, implying that the sets D_a are a.s. empty for all $T < \infty$. Since a.s.

$$\liminf_{\epsilon\to 0} \frac{\mu_\infty^W(B(0,\epsilon))}{\epsilon^2} = 0$$

(see [29, Theorem 6.8]), taking $k \uparrow \infty$ and $a \downarrow C_d$ completes the proof of (1.16) and hence of Theorem 1.6, subject only to (9.2).

The first step in the proof of (9.2) is the following simple lemma (see [22] for the definition and properties of Bessel processes).

LEMMA 9.1. *Let $Z = \int_0^T U_s^{-2} ds$ with $\{U_s : s \in [0,T]\}$ the Bessel process of index $d' = d/2 - 1 > 0$, starting at $U_0 = u \in (0,k]$. Then, for any $\alpha \in (0, d']$, $b > 1$, there exist $c = c(b, T, d', k) < \infty$ such that*

$$(9.7) \qquad \mathbb{E}^u_{(d')}\big(e^{(d'^2/2-\alpha^2/2)Z}\mathbf{1}_{\inf_{s\in[0,T]} U_s \le v}\big) \le c v^{2\alpha/b} u^{-(d'-\alpha)-2\alpha/b} \, .$$

Proof: Let $P^u_{(\nu)}(\cdot)$ denote the law of the Bessel process $\{U_s : s \in [0,T]\}$ of index $\nu > 0$ starting at $U_0 = u$. Recall that for any index $\nu > 0$,

$$dU_s = (\nu + 1/2)\frac{ds}{U_s} + dB_s \, , \quad U_0 = u > 0 \, ,$$

where B_s is a one dimensional Brownian motion. In particular, $dP^u_{(d')}/dP^u_{(\alpha)}$ exists for any $d' \ge \alpha > 0$ and is given by the Girsanov transformation as (see [22, Pg. 419]),

$$\frac{dP^u_{(d')}}{dP^u_{(\alpha)}} = (\frac{U_T}{u})^{d'-\alpha} e^{-(d'^2/2-\alpha^2/2)\int_0^T U_s^{-2} ds} \, .$$

In particular, by Hölder's inequality, for $q = b/(b-1)$,

$$\begin{aligned}
\mathbb{E}^u_{(d')}\big(e^{(d'^2/2-\alpha^2/2)Z}\mathbf{1}_{\inf_{s\in[0,T]} U_s \le v}\big) &= u^{\alpha-d'}\mathbb{E}^u_{(\alpha)}\big(U_T^{d'-\alpha}\mathbf{1}_{\inf_{s\in[0,T]} U_s \le v}\big) \\
&\le u^{\alpha-d'} P^u_{(\alpha)}(\inf_{s\ge 0} U_s \le v)^{1/b}\mathbb{E}^u_{(\alpha)}(U_T^{q(d'-\alpha)})^{1/q} \\
&\le u^{\alpha-d'}(\frac{v}{u})^{2\alpha/b}\mathbb{E}^k_{(\alpha)}(U_T^{q(d'-\alpha)})^{1/q}
\end{aligned}$$

where the last inequality follows using the fact that the Bessel process of index α has scale function $-x^{-2\alpha}$ [22, Pg. 415], and for the right-most expectation we used a simple comparison argument. A further comparison argument shows that we can take $c = \mathbb{E}^k_{(d')}((1 \vee U_T)^{qd'})^{1/q} < \infty$. $\square$

The next step in proving (9.2) is to establish the following consequence of Lemma 9.1.

LEMMA 9.2. *For any $T < \infty$, $b > 1$, $k > 1$, there exists $c < \infty$ such that for any $a > 0$, $\alpha \in (0, d']$, $\eta > 0$, $\delta = \eta^{1-b^{-1}}$, $|x| \in (0, k]$,*

$$\mathbf{P}\Big(\inf_{\epsilon \in [\eta,\delta]} \frac{\mu_T^W(B(x, b\epsilon))}{\epsilon^2} \geq \frac{a}{b}\Big) \leq c\eta^{ab^{-4}(d'^2-\alpha^2)+2\alpha/b}|x|^{-(d'-\alpha)-2\alpha/b} . \tag{9.8}$$

Proof: Fix $T, a, b, k, \alpha, \eta, \delta$ and x as in the statement of the lemma. Observe that $U_s = |W_s - x|$ is a Bessel process of index d', starting at $U_0 = |x| \in (0, k]$. Clearly

$$\{\mu_T^W(B(x, v)) > 0\} = \{\inf_{s \in [0,T]} U_s < v\} \tag{9.9}$$

Setting $Z = \int_0^T U_s^{-2} ds$, also

$$\begin{aligned} b^2 Z &= \int_0^T \int_{b^{-1}U_s}^\infty \frac{2d\epsilon}{\epsilon^3}\, ds = \int_0^T \int_0^\infty \mathbf{1}_{\{|W_s - x| \leq b\epsilon\}} \frac{2d\epsilon}{\epsilon^3}\, ds \\ &= \int_0^\infty \frac{2d\epsilon}{\epsilon^3} \mu_T^W(B(x, b\epsilon)) \geq \int_\eta^\delta \frac{2d\epsilon}{\epsilon^3} \mu_T^W(B(x, b\epsilon)) . \end{aligned} \tag{9.10}$$

If

$$\inf_{\epsilon \in [\eta,\delta]} \epsilon^{-2} \mu_T^W(B(x, b\epsilon)) \geq \frac{a}{b}$$

then $\mu_T^W(B(x, b\eta)) > 0$ and

$$\int_\eta^\delta \frac{2d\epsilon}{\epsilon^3} \mu_T^W(B(x, b\epsilon)) \geq \frac{a}{b} \int_\eta^\delta \frac{2d\epsilon}{\epsilon} = -2ab^{-2} \log \eta .$$

Thus, for $v = b\eta$ and $\lambda = (d'^2 - \alpha^2)/2 \geq 0$, by (9.9), (9.10) and Chebycheff's inequality,

$$\mathbf{P}\Big(\inf_{\epsilon \in [\eta,\delta]} \frac{\mu_T^W(B(x, b\epsilon))}{\epsilon^2} \geq \frac{a}{b}\Big) \leq P^{|x|}_{(d')}(Z \geq -2ab^{-4} \log \eta, \inf_{s \in [0,T]} U_s \leq v)$$

$$\leq \eta^{2\lambda ab^{-4}} \mathbb{E}^{|x|}_{(d')}[e^{\lambda Z} \mathbf{1}_{\inf_{s\in[0,T]} U_s \leq v}]$$

We thus obtain (9.8) by applying Lemma 9.1. □

We now return to complete the proof of (9.2). For $b > 1$ and $\alpha \in (0, d']$ let

$$f_a(b,\alpha) = ab^{-4}(d'^2 - \alpha^2) - d + 2\alpha/b\ .$$

By Lemma 9.2, for some $c, c_1, c_2 < \infty$ independent of n,

$$\begin{aligned}
\mathbb{E}|\mathcal{H}_n| &= \sum_{j=1}^{K_n} \mathbf{P}\big(\inf_{\epsilon\in[\eta_n,\delta_n]} \frac{\mu_T^W(B(x_j, b\epsilon))}{\epsilon^2} \geq \frac{a}{b}\big) \\
&\leq c\eta_n^{ab^{-4}(d'^2-\alpha^2)+2\alpha/b} \sum_{j=1}^{K_n} |x_j|^{-(d'-\alpha)-2\alpha/b} \\
&\leq c_1 \eta_n^{ab^{-4}(d'^2-\alpha^2)+2\alpha/b-d}\big(1 + \int_{\{|x|\leq k\}} |x|^{-(d'-\alpha)-2\alpha/b}\, dx\big) \\
&\leq c_2 \eta_n^{f_a(b,\alpha)}, \qquad (9.11)
\end{aligned}$$

using $(d'-\alpha) + 2\alpha/b < d' + \alpha \leq d-1$. inequality).

Setting $\alpha = d' - \theta$ for $\theta \in [0, d')$, in which case $d'^2 - \alpha^2 = \theta(d-2-\theta)$, we see that

$$f_a(b,\alpha) = ab^{-4}\theta(d-2-\theta) - d(1-b^{-1}) - (2\theta+2)/b. \qquad (9.12)$$

Observe that $I_d(a)$, defined in (1.12), can also be written as

$$(\max\{0, a(d-2)-2\})^2/4a\,,$$

whence

$$I_d(a) = \sup_{0\leq\theta<(d-2)/2} \{a\theta(d-2-\theta) - 2\theta\}\,, \qquad (9.13)$$

and the supremum in (9.13) is attained at $\theta = \max\{0, (d-2)/2 - a^{-1}\}$. Comparing this with (9.12) we see that

$$\lim_{b\downarrow 1} \sup_{\alpha\in(0,d']} f_a(b,\alpha) = I_d(a) - 2 \qquad (9.14)$$

which completes the proof of (9.2) and hence of Theorem 1.6. $\square$

Some unsolved problems:

1. Determine exactly the dimension appearing in (1.14) and the precise asymptotics in (1.16).

2. Does the set considered in (1.14) have equal Hausdorff and packing dimensions?

3. By arguments similar to those in the proof of (1.16), we can show that there exist non-random constants $\tilde{c}_d > 0$, $\tilde{C}_d < \infty$ such that

$$\tilde{c}_d \le \inf_{t\in[0,1]} \limsup_{\epsilon\to 0} \frac{\mu_\infty^W(B(W_t,\epsilon))}{\epsilon^2} \le \tilde{C}_d \quad \text{a.s.} \tag{9.15}$$

More precisely, the upper bound here is proved just like the lower bound in (1.16), while the lower bound can be inferred from Perkins [17] or from a branching process argument. As in (1.16), it is an open problem to determine the optimal constants in (9.15).

Acknowledgments. We thank N.-R. Shieh and S. J. Taylor for useful discussions and for providing us with a preprint of [24], and J. Pitman for a useful discussion of the Ciesielski-Taylor identity. We are indebted to MSRI, to the organizers of the program on stochastic analysis there, and in particular to S. Evans and R. Williams, for making our collaboration possible.

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