## Numerical Integration of 3D Reaction-Diffusion Equations (Difference methods for parabolic PDEs, boundary value problem)

### **Reaction-Diffusion Problem:**

Calcium : 
$$\frac{\partial [Ca^{2+}]}{\partial t} = D_{Ca} \nabla^2 [Ca^{2+}] + \sum_j R_j + \sum_m a_m \delta(x - x_m)$$
(1)

Buffer : 
$$\frac{\partial[B_j]}{\partial t} = D_{B_j} \nabla^2[B_j] + R_j$$
 (2)

Reaction term (  $Ca + B_j \rightleftharpoons CaB_j$  ) :

$$R_j = -k_j^+[B_j][Ca^{2+}] + k_j^-([B_j]_{total} - [B_j])$$
(3)

- Diffusion/heat equation in one dimension
  - Explicit and implicit difference schemes
  - Stability analysis
  - Non-uniform grid
- Three dimensions: Alternating Direction Implicit (ADI) methods
- Non-homogeneous diffusion equation: dealing with the reaction term

### Diffusion/heat equation in one spatial dimension

$$\frac{\partial f}{\partial t} = D \frac{\partial^2 f}{\partial x^2} \tag{4}$$

Discretization of coordinates:  $(x, t) \rightarrow (x_j, t_n)$  $x_j = x_0 + j \Delta x$ ;  $t_n = t_0 + n \Delta t$ ;  $f(x_j, t_n) \equiv f_j^n$ 

Finite difference approximation for 1st derivative:

$$\frac{f_j^{n+1} - f_j^n}{\Delta t} = \frac{\partial f_j^{n+1/2}}{\partial t} + O(\Delta t^2) = \frac{\partial f_j^n}{\partial t} + O(\Delta t) = \frac{\partial f_j^{n+1}}{\partial t} + O(\Delta t)$$
(5)

Finite difference approximation for 2nd derivative:

$$\frac{(f_{j+1}^n - f_j^n)/\Delta x - (f_j^n - f_{j-1}^n)/\Delta x}{\Delta x} = \frac{f_{j-1}^n - 2f_j^n + f_{j+1}^n}{\Delta x^2} = \frac{\partial^2 f_j^n}{\partial x^2} + O(\Delta x^2) \quad (6)$$

1. Euler (explicit) scheme:  $f_j^{n+1} - f_j^n = \nu (f_{j-1}^n - 2f_j^n + f_{j+1}^n) \equiv \nu \, \delta_x^2 f_j^n$ where  $\nu \equiv D\Delta t / \Delta x^2$ 

- Given  $f_j^n$ , values at next time step  $f_j^{n+1}$  are computed directly
- Stable for  $\nu < 1/2$
- Truncation error:  $T = O(\Delta t) + O(\Delta x^2)$

2. Fully implicit scheme:  $f_j^{n+1} - f_j^n = \nu \left( f_{j-1}^{n+1} - 2f_j^{n+1} + f_{j+1}^{n+1} \right) \equiv \nu \, \delta_x^2 f_j^{n+1}$  ("backward" Euler method)

- Tridiagonal system of linear equations must be solved at each time step
- Unconditionally stable
- Truncation error:  $T = O(\Delta t) + O(\Delta x^2)$

# 3. Crank-Nicholson (implicit) scheme: $f_j^{n+1} - f_j^n = \nu/2\{\delta_x^2 f_j^n + \delta_x^2 f_j^{n+1}\}$ symmetric representation: $(1 - \frac{\nu}{2}\delta_x^2)f_j^{n+1} = (1 + \frac{\nu}{2}\delta_x^2)f_j^n$ (7)

- Truncation error:  $T = O(\Delta t^2) + O(\Delta x^2)$
- Unconditionally stable

Solving the fully implicit scheme: tridiagonal system

$$\frac{\partial f}{\partial t} = D \frac{\partial^2 f}{\partial x^2} \qquad \Rightarrow \qquad f_j^{n+1} - f_j^n = \nu \left( f_{j-1}^{n+1} - 2f_j^{n+1} + f_{j+1}^{n+1} \right) \tag{8}$$

• Assume Dirichlet boundary conditions:  $f_0 = f_N = 0$ 

$$\begin{pmatrix} 1+2\nu & -\nu & 0 & 0 & \cdot & & \\ -\nu & 1+2\nu & -\nu & 0 & \cdot & 0 & \\ 0 & -\nu & 1+2\nu & -\nu & \cdot & & \\ 0 & 0 & -\nu & 1+2\nu & \cdot & & \\ & & & \cdot & 1+2\nu & -\nu & 0 \\ & 0 & & & -\nu & 1+2\nu & -\nu \\ & & & & 0 & -\nu & 1+2\nu \end{pmatrix} \begin{pmatrix} f_0^{n+1} \\ f_1^{n+1} \\ f_3^{n} \\ f_4^{n} \\ \cdot \\ f_{N-2}^{n+1} \\ f_N^{n+1} \\ f_N^{n+1} \\ f_N^{n+1} \end{pmatrix} = \begin{pmatrix} f_0^n \\ f_1^n \\ f_3^n \\ f_4^n \\ \cdot \\ f_{N-2}^n \\ f_N^{n-1} \\ f_N^n \end{pmatrix}$$

$$L_j f_{j-1} + C_j f_j + R_j f_{j+1} = H_j \qquad j = 0..N$$

$$L_0 = R_N = 0, \ L_j = R_j = -\nu, \ C_j = 1 + 2\nu$$

#### Solve by Gaussian elimination:

1. Forward elimination:  $R_0 = R_0/C_0$ ,  $H_0 = H_0/C_0$  (9)

$$C_j = C_j - L_j R_{j-1}, \quad j = 1..N$$
 (10)

$$H_j = (H_j - L_j H_{j-1})/C_j, \quad j = 1..N$$
 (11)

$$R_j = R_j / C_j, \quad j = 1..N$$
 (12)

2. Backward substitution:  $f_N = H_N$  (13)

$$f_j = H_j - R_j f_{j+1}, \quad j = N - 1 \dots 0.$$
 (14)

• Operations per time step: 5N multiply/divide + 3N add/subtract

### Stability Analysis: Fourier Method

• Difference scheme is linear (as well as the PDEs), so the error satisfies the same equation

• Substitute into equations the solution (for the error instability) of the form  $f_j^n = \lambda^n \exp(i x_j k)$  ( where  $k = \pi m/(N\Delta x)$ , m = 0..N )

- Method is stable if  $\lambda < 1$
- 1. Forward Euler (explicit scheme):  $f_j^{n+1} f_j^n = \nu_x (f_{j-1}^n 2f_j^n + f_{j+1}^n)$

$$(\lambda^{n+1} - \lambda^n) \exp(i x_j k) = \nu_x \lambda^n \exp(i x_j k) \left[ \exp(-i \Delta x k) - 2 + \exp(i \Delta x k) \right]$$
$$\lambda - 1 = \nu_x \left[ \exp(-i \Delta x k) - 2 + \exp(i \Delta x k) \right]$$
$$\lambda = 1 - 4\nu_x \sin^2(\Delta x k/2)$$
$$|\lambda| < 1 \rightarrow \nu_x < 1/2$$
(15)

2. Backward Euler (implicit scheme):  $f_j^{n+1} - f_j^n = \nu_x (f_{j-1}^{n+1} - 2f_j^{n+1} + f_{j+1}^{n+1})$ 

$$(\lambda^{n+1} - \lambda^n) \exp(i x_j k) = \nu_x \lambda^{n+1} \exp(i x_j k) \left[\exp(-i \Delta x k) - 2 + \exp(i \Delta x k)\right]$$
$$\lambda - 1 = \nu_x \lambda \left[\exp(-i \Delta x k) - 2 + \exp(i \Delta x k)\right]$$
$$\lambda = \frac{1}{1 + 4\nu_x \sin^2(\Delta x k/2)} < 1 \quad \text{(always stable)} \tag{16}$$

# **Non-Uniform Grid:** $x_{i+1} - x_i = \Delta x_i$

Finite difference approximation for the 2nd derivative:

$$\frac{(f_{j+1}^n - f_j^n)/\Delta x_j - (f_j^n - f_{j-1}^n)/\Delta x_{j-1}}{(\Delta x_{j-1} + \Delta x_j)/2} = \frac{\partial^2 f_j^n}{\partial x^2} + O(\Delta x_{j,j-1}^2) + O(\Delta x_j - \Delta x_{j-1})$$
(17)

• To maintain accuracy, set  $\Delta x_j = \Delta x_{j-1}(1 \pm \epsilon), \ \epsilon \ll 1$ 

Modified difference scheme:

 $f_{j}^{n+1} - f_{j}^{n} = \nu \left( f_{j-1}^{n+1} - 2f_{j}^{n+1} + f_{j+1}^{n+1} \right) \quad \Rightarrow \quad f_{j}^{n+1} - f_{j}^{n} = \nu_{j}^{-} f_{j-1}^{n+1} + \nu_{j}^{0} f_{j}^{n+1} + \nu_{j}^{+} f_{j+1}^{n+1} \tag{18}$ 

where

$$\nu_{j}^{-} = \frac{2D\Delta t}{\Delta x_{j-1}(\Delta x_{j-1} + \Delta x_{j})}$$
$$\nu_{j}^{0} = -\frac{2D\Delta t}{\Delta x_{j-1}\Delta x_{j}}$$
$$\nu_{j}^{+} = \frac{2D\Delta t}{\Delta x_{j}(\Delta x_{j-1} + \Delta x_{j})}$$

Minimal spacing decreases exponentially with increasing N:

$$\Delta x_j = (1+\epsilon)\Delta x_{j-1} = (1+\epsilon)^j \Delta x_0 \equiv \gamma^j \Delta x_0$$
$$\Delta x_0 = \frac{\Delta x_N}{\gamma^n}$$

Flux boundary conditions: 
$$\frac{\partial f_{0,N}}{\partial n} + b f_{0,N} = c$$
 (19)

Consider field values at (virtual) points outside the boundary:  $f_{-1} = f(x_0 - \Delta x), f_{N+1} = f(x_N + \Delta x)$ 

$$\frac{\partial f_0}{\partial x} \Rightarrow \frac{f_1 - f_{-1}}{2\Delta x} + O(\Delta x^2)$$
  
$$\frac{\partial f_0}{\partial x} + bf_0 = c \Rightarrow f_1 - f_{-1} + 2\Delta x (bf_0 - c) = 0$$
  
$$f_{-1} = f_1 + 2\Delta x (bf_0 - c)$$
(20)

Difference equation at point  $x_0$ :

$$f_0^{n+1} - f_0^n = \nu \left( f_{-1}^{n+1} - 2f_0^{n+1} + f_1^{n+1} \right) = \nu \left( f_1^{n+1} + 2\Delta x (bf_0^{n+1} - c) - 2f_0^{n+1} + f_1^{n+1} \right) = 2\nu \left( f_1^{n+1} + (b\Delta x - 1)f_0^{n+1} - c\Delta x \right) \Rightarrow \left[ 2\nu (1 - b\Delta x) + 1 \right] f_0^{n+1} - 2\nu f_1^{n+1} = f_0^n - 2\nu c\Delta x$$
(21)

Difference equation at point  $x_N$ :

$$[2\nu(1 - b\Delta x) + 1]f_N^{n+1} - 2\nu f_{N-1}^{n+1} = f_N^n - 2\nu c\Delta x$$
(22)

$$\begin{pmatrix} 1+2\nu(1-b\Delta x) & -2\nu & 0 & \cdot \\ & -\nu & 1+2\nu & -\nu & \cdot & 0 \\ 0 & -\nu & 1+2\nu & \cdot & & \\ & \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ & & & \cdot & 1+2\nu & -\nu & 0 \\ 0 & & \cdot & -\nu & 1+2\nu & -\nu \\ & & 0 & -2\nu & 1+2\nu(1-b\Delta x) \end{pmatrix} \begin{pmatrix} f_0^{n+1} \\ f_1^{n+1} \\ f_3^{n+1} \\ \vdots \\ f_{N-1}^{n+1} \\ f_N^{n+1} \end{pmatrix}$$
$$= \begin{pmatrix} f_0^n - 2\nu c\Delta x \\ f_1^n \\ f_3^n \\ \cdot \\ f_{N-2}^n \\ f_{N-1}^n \\ f_N^n - 2\nu c\Delta x \end{pmatrix}$$

### Diffusion/heat equation in three spatial dimensions

$$\frac{\partial f}{\partial t} = D \left[ \frac{\partial^2 f}{\partial x^2} + \frac{\partial^2 f}{\partial y^2} + \frac{\partial^2 f}{\partial z^2} \right]$$
(23)

1. Euler (explicit) scheme:  $f^{n+1} - f^n = \left[\nu_x \, \delta_x^2 + \nu_y \, \delta_y^2 + \nu_z \, \delta_z^2\right] f^n$ where  $f^n \equiv f_{ijk}^n \equiv f(x_i, y_j, z_k; t^n)$  $\nu_x \equiv D\Delta t/\Delta x^2, \, \nu_y \equiv D\Delta t/\Delta y^2, \, \nu_z \equiv D\Delta t/\Delta z^2$  $\delta_x^2 f^n \equiv f_{i-1jk}^n - 2f_{ijk}^n + f_{i+1jk}^n, \, \delta_y^2 f^n \equiv f_{ij-1k}^n - 2f_{ijk}^n + f_{ij+1k}^n$  $\delta_z^2 f^n \equiv f_{ijk-1}^n - 2f_{ijk}^n + f_{ijk+1}^n$ 

- Stability condition:  $\nu_x + \nu_y + \nu_z \le 1/2$
- 2. Implicit schemes: Alternating Direction Implicit (ADI) Methods

Backward Euler  $(f^{n+1} - f^n = [\nu_x \delta_x^2 + \nu_y \delta_y^2 + \nu_z \delta_z^2] f^{n+1})$  no longer practical: requires solution of a (non-tridiagonal) system of  $N^3$  equations.

Locally One-Dimensional (LOD) method

3D analog of Crank-Nicholson scheme:

$$\left(1 - \frac{\nu_x}{2}\delta_x^2 - \frac{\nu_y}{2}\delta_y^2 - \frac{\nu_z}{2}\delta_z^2\right)f^{n+1} = \left(1 + \frac{\nu_x}{2}\delta_x^2 + \frac{\nu_y}{2}\delta_y^2 + \frac{\nu_z}{2}\delta_z^2\right)f^n \tag{24}$$

• Factorization (introduces additional terms of order  $O(\Delta t^2)$ )

$$\left(1 - \frac{\nu_x}{2}\delta_x^2\right)\left(1 - \frac{\nu_y}{2}\delta_y^2\right)\left(1 - \frac{\nu_z}{2}\delta_z^2\right)f^{n+1} = \left(1 + \frac{\nu_x}{2}\delta_x^2\right)\left(1 + \frac{\nu_y}{2}\delta_y^2\right)\left(1 + \frac{\nu_z}{2}\delta_z^2\right)f^n \quad (25)$$

• To solve Eq. (25), break down each iteration into several steps:

$$(1 - \frac{\nu_x}{2}\delta_x^2)f^{n*} = (1 + \frac{\nu_x}{2}\delta_x^2)f^n (1 - \frac{\nu_y}{2}\delta_y^2)f^{n**} = (1 + \frac{\nu_y}{2}\delta_y^2)f^{n*} (1 - \frac{\nu_z}{2}\delta_z^2)f^{n+1} = (1 + \frac{\nu_z}{2}\delta_z^2)f^{n**}$$
 (26)

- Each time step three tridiagonal systems are solved for all points.
- Truncation error:  $T = O(\Delta t^2) + O(\Delta x^2) + O(\Delta y^2) + O(\Delta z^2)$
- Always stable

Problem: equations for intermediate values  $f^{n*}$ ,  $f^{n**}$  are not consistent with the full differential equation: additional care required dealing with boundary conditions

### **Douglas-Gunn ADI Method**

Another "deformation" of the Crank-Nicholson method:

$$(1 - \frac{A_x}{2})(1 - \frac{A_y}{2})(1 - \frac{A_z}{2})f^{n+1} = \left[(1 + \frac{A_x}{2})(1 + \frac{A_y}{2})(1 + \frac{A_z}{2}) - \frac{A_x A_y A_z}{4}\right]f^n \quad (27)$$
  
where  $A_x \equiv \nu_x \delta_x^2 = D\frac{\Delta t}{\Delta x^2} \delta_x^2, \ A_y \equiv \nu_y \delta_y^2, \ A_z \equiv \nu_z \delta_z^2$ 

Multi-step implementation - Douglas-Gunn method

$$(1 - \frac{A_x}{2})f^{n*} = (1 + \frac{A_x}{2} + A_y + A_z)f^n$$
(28)

$$(1 - \frac{A_y}{2})f^{n**} = (1 + \frac{A_x}{2} + \frac{A_y}{2} + A_z)f^n + \frac{A_x}{2}f^{n*}$$
(29)

$$(1 - \frac{A_z}{2})f^{n+1} = (1 + \frac{A_x}{2} + \frac{A_y}{2} + \frac{A_z}{2})f^n + \frac{A_x}{2}f^{n*} + \frac{A_y}{2}f^{n**}$$
(30)

# • Each equation is a valid approximation of the full diffusion equation - no modification needed for boundary conditions

Simplify by subtracting Eq.(28) from Eq.(29), and Eq.(29) from Eq.(30):

$$(1 - \frac{A_x}{2})f^{n*} = (1 + \frac{A_x}{2} + A_y + A_z)f^n$$
(31)

$$(1 - \frac{A_y}{2})f^{n**} = f^{n*} - \frac{A_y}{2}f^n$$
(32)

$$(1 - \frac{A_z}{2})f^{n+1} = f^{n**} - \frac{A_z}{2}f^n$$
(33)

- Truncation error:  $T = O(\Delta t^2) + O(\Delta x^2) + O(\Delta y^2) + O(\Delta z^2)$
- Always stable

### Heterogeneous Diffusion/heat equation in 1D

$$\frac{\partial f}{\partial t} = D \frac{\partial^2 f}{\partial x^2} + H(f, t) \tag{34}$$

**Crank-Nicholson scheme:** 

$$f_j^{n+1} - f_j^n = \nu/2\{\delta_x^2 f_j^n + \delta_x^2 f_j^{n+1}\} + \frac{\Delta t}{2}(H(f^{n+1}, t^{n+1}) + H(f^n, t^n))$$
(35)

Symmetric form:

$$\left(1 - \frac{\nu}{2}\delta_x^2\right)f_j^{n+1} = \left(1 + \frac{\nu}{2}\delta_x^2\right)f_j^n + \frac{\Delta t}{2}\left(H(f^{n+1}, t^{n+1}) + H(f^n, t^n)\right)$$
(36)

• For non-linear H(f), this scheme can be solved using predictor/corrector or iteration methods, or one has to give up  $O(\Delta t^2)$  accuracy:  $H(f^{n+1}) = H(f^n) + H'(f^n)(f^{n+1} - f^n)/\Delta t + O(\Delta t)$ 

• Without loss of accuracy, one can substitute  $H^n \equiv H(f^n, t^n) \rightarrow H(f^n, t^{n+1/2})$ ,  $H^{n+1} \equiv H(f^{n+1}, t^{n+1}) \rightarrow H(f^{n+1}, t^{n+1/2})$ . The following is an equivalent scheme:

$$(1 - \frac{\nu}{2}\delta_x^2)f_j^{n+1} = (1 + \frac{\nu}{2}\delta_x^2)f_j^n + \frac{\Delta t}{2}(H(f^{n+1}, t^{n+\frac{1}{2}}) + H(f^n, t^{n+\frac{1}{2}}))$$
(37)

### Heterogeneous diffusion/heat equation in 3D

$$\frac{\partial f}{\partial t} = D \left[ \frac{\partial^2 f}{\partial x^2} + \frac{\partial^2 f}{\partial y^2} + \frac{\partial^2 f}{\partial z^2} \right] + H(f, t)$$
(38)

Douglas-Gunn method for heterogeneous equation:

$$(1 - \frac{A_x}{2})f^{n*} = (1 + \frac{A_x}{2} + A_y + A_z)f^n + \Delta t \left(\alpha H(f^n) + \beta H(f^{n*})\right)$$
(39)

$$(1 - \frac{A_y}{2})f^{n**} = (1 + \frac{A_x}{2} + \frac{A_y}{2} + A_z)f^n + \frac{A_x}{2}f^{n*} + \Delta t\left(\gamma H(f^n) + \delta H(f^{n**})\right)$$
(40)

$$(1 - \frac{A_z}{2})f^{n+1} = (1 + \frac{A_x}{2} + \frac{A_y}{2} + \frac{A_z}{2})f^n + \frac{A_x}{2}f^{n*} + \frac{A_y}{2}f^{n**} + \Delta t\left(\zeta H(f^n) + \eta H(f^{n+1})\right)$$
(41)

where  $\alpha + \beta = 1$ ,  $\gamma + \delta = 1$ ,  $\zeta + \eta = 1$ . Choice of coefficients consistent with second order accuracy:  $\alpha = \gamma = 1$ ,  $\beta = \delta = 0$ ,  $\zeta = \eta = 1/2$ .

Subtract Eq.(39) from Eq.(40), and Eq.(40) from Eq.(41):

$$(1 - \frac{A_x}{2})f^{n*} = (1 + \frac{A_x}{2} + A_y + A_z)f^n + \Delta t H^n$$
(42)

$$(1 - \frac{A_y}{2})f^{n**} = f^{n*} - \frac{A_y}{2}f^n$$
(43)

$$(1 - \frac{A_z}{2})f^{n+1} = f^{n**} - \frac{A_z}{2}f^n + \frac{\Delta t}{2}(H^{n+1} - H^n)$$
(44)

where  $H^n \equiv H(f^n, t^n), H^{n+1} \equiv H(f^{n+1}, t^{n+1}).$ 

Apply operator  $(1 - \frac{A_x}{2})(1 - \frac{A_y}{2})$  from the left to recover the difference scheme:

$$(1 - \frac{A_x}{2})(1 - \frac{A_y}{2})(1 - \frac{A_z}{2})f^{n+1} = \left[(1 + \frac{A_x}{2})(1 + \frac{A_y}{2})(1 + \frac{A_z}{2}) - \frac{A_xA_yA_z}{4}\right]f^n + \frac{\Delta t}{2}(H^n + H^{n+1}) - \frac{\Delta t}{2}(A_x + A_y - \frac{A_xA_y}{2})(H^{n+1} - H^n)$$
(45)

• Without loss of accuracy, one can substitute  $H^n \equiv H(f^n, t^n) \to H(f^n, t^{n+1/2})$ ,  $H^{n+1} \equiv H(f^{n+1}, t^{n+1}) \to H(f^{n+1}, t^{n+1/2})$ .

## Reaction-diffusion with one buffer (Hines method)

Calcium : 
$$\frac{\partial [Ca]}{\partial t} = D_{Ca} \nabla^2 [Ca] + R + \sum_m a_k \delta(x - x_m)$$
 (46)

Buffer : 
$$\frac{\partial[B]}{\partial t} = D_B \nabla^2[B] + R$$
 (47)

Reaction term : 
$$R = -k^+[B][Ca] + k^-([B]_{total} - [B])$$
 (48)

• Difference scheme for source term:  $H_{source}(x_i, y_j, z_k) = \sum c_m \delta^i_{i_m} \delta^j_{j_m} \delta^k_{k_m}$ 

Heterogeneity for [Ca]: 
$$H_{Ca}([Ca], t) = H_{source} + R([Ca], [B])$$
  
Heterogeneity for [B]:  $H_B([B], t) = R([Ca], [B])$  (49)

• **Problem:** to solve for  $[Ca]^{n+1}$ , knowledge of  $[B]^{n+1}$  is required; to solve for  $[B]^{n+1}$ , knowledge of  $[Ca]^{n+1}$  is required.

• Solution: compute [Ca] and [B] on time grids staggered by  $\Delta t/2$ 

$$(1 - \frac{A_x}{2})[Ca]^{n*} = (1 + \frac{A_x}{2} + A_y + A_z)[Ca]^n + \Delta t H_{Ca}([Ca]^n, [B]^{n+\frac{1}{2}}))$$
  

$$(1 - \frac{A_y}{2})[Ca]^{n**} = [Ca]^{n*} - \frac{A_y}{2}[Ca]^n$$
  

$$(1 - \frac{A_z}{2})[Ca]^{n+1} = [Ca]^{n**} - \frac{A_z}{2}[Ca]^n + \frac{\Delta t}{2}(H_{Ca}([Ca]^{n+1}, [B]^{n+\frac{1}{2}}) - H_{Ca}([Ca]^n, [B]^{n+\frac{1}{2}}))$$

$$(1 - \frac{A_x}{2})[B]^{n - \frac{1}{2}*} = (1 + \frac{A_x}{2} + A_y + A_z)[B]^{n - \frac{1}{2}} + \Delta t H_B([B]^{n - \frac{1}{2}}, [Ca]^n) (1 - \frac{A_y}{2})[B]^{n - \frac{1}{2}**} = [B]^{n - \frac{1}{2}*} - \frac{A_y}{2}[B]^{n - \frac{1}{2}} (1 - \frac{A_z}{2})[B]^{n + \frac{1}{2}} = [B]^{n - \frac{1}{2}**} - \frac{A_z}{2}[B]^{n - \frac{1}{2}} + \frac{\Delta t}{2}(H_B([B]^{n + \frac{1}{2}}, [Ca]^n) - H_B([B]^{n - \frac{1}{2}}, [Ca]^n))$$

### Fully implicit difference scheme in 3D with heterogeneity

Fully implicit ADI method (Douglas-Rachford):

$$(1 - A_x)f^{n*} = (1 + A_y + A_z)f^n + H^n\Delta t$$
(50)

$$(1 - A_y)f^{n**} = (1 + A_z)f^n + A_x f^{n*} + H^n \Delta t$$
(51)

$$(1 - A_z)f^{n+1} = f^n + A_x f^{n*} + A_y f^{n**} + H^{n+1}\Delta t$$
(52)

• Each equation is a valid approximation of the full diffusion equation - no modification needed for boundary conditions

Agrees with the following "deformation" of the fully implicit difference scheme:

$$(1 - A_x)(1 - A_y)(1 - A_z)f^{n+1} = (1 + A_xA_y + A_yA_z + A_xA_z - A_xA_yA_z)f^n + H^{n+1}\Delta t + A_xA_y(H^{n+1} - H^n)\Delta t$$
(53)

where  $A_x \equiv \nu_x \delta_x^2 = D \frac{\Delta t}{\Delta x^2} \delta_x^2$ ,  $A_y \equiv \nu_y \delta_y^2$ ,  $A_z \equiv \nu_z \delta_z^2$ 

Simplify by subtracting Eq.(50) from Eq.(51), and Eq.(51) from Eq.(52):

$$(1 - A_x)f^{n*} = (1 + A_y + A_z)f^n + H^n \Delta t$$
(54)

$$(1 - A_y)f^{n**} = -A_y f^n + f^{n*}$$
(55)

$$(1 - A_z)f^{n+1} = -A_z f^n + f^{n**} + (H^{n+1} - H^n)\Delta t$$
(56)

• Truncation error:  $T = O(\Delta t) + O(\Delta x^2) + O(\Delta y^2) + O(\Delta z^2)$ 

• Always stable

# Heterogeneous diffusion/heat equation in 2D

$$\frac{\partial f}{\partial t} = D \left[ \frac{\partial^2 f}{\partial x^2} + \frac{\partial^2 f}{\partial y^2} \right] + H(f, t)$$
(57)

### 1. Douglas-Gunn/Crank-Nicholson method:

$$(1 - \frac{A_x}{2})f^{n*} = (1 + \frac{A_y}{2})f^n + \frac{\Delta t}{2}H(f^n, t^n)$$
(58)

$$(1 - \frac{A_y}{2})f^{n+1} = (1 + \frac{A_x}{2})f^{n*} + \frac{\Delta t}{2}H(f^{n+1}, t^{n+1})$$
(59)

Corresponding difference scheme:

$$(1 - \frac{A_x}{2})(1 - \frac{A_y}{2})f^{n+1} = (1 + \frac{A_x}{2})(1 + \frac{A_y}{2})f^n + \frac{\Delta t}{2}(H^n + H^{n+1}) + \frac{\Delta t}{4}A_x(H^n - H^{n+1})$$
(60)

where  $A_x \equiv \nu_x \delta_x^2 = D \frac{\Delta t}{\Delta x^2} \delta_x^2$ ,  $A_y \equiv \nu_y \delta_y^2$ 

### 2. Fully implicit method:

$$(1 - A_x)f^{n*} = (1 + A_y)f^n + \Delta t H(f^n, t^n)$$
(61)

$$(1 - A_y)f^{n+1} = f^n + A_x f^{n*} + \Delta t H(f^{n+1}, t^{n+1})$$
(62)

Corresponding "deformation" of the fully implicit difference scheme:

$$(1 - A_x)(1 - A_y)f^{n+1} = (1 + A_xA_y)f^n + H^{n+1}\Delta t + A_x(H^n - H^{n+1})\Delta t$$
(63)